New VLT observations of the *Fermi* pulsar PSR J1048–5832

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ABSTRACT

PSR J1048–5832 is a Vela-like (P=123.6 ms; τ ∼ 20.3 kyr) γ-ray pulsar detected by Fermi, at a distance of ∼ 2.7 kpc and with a rotational energy loss rate $E_{SD} \sim 2 \times 10^{36}$ erg s$^{-1}$. The PSR J1048–5832 field has been observed with the Very Large Telescope (VLT) in the V and R bands. We used these data to determine the colour of the object detected closest to the *Chandra* position (Star D) and confirm that it is not associated with the pulsar. For the estimated extinction along the line of sight, inferred from a re-analysis of the *Chandra* and XMM-Newton spectra, the fluxes of Star D ($V \sim 26.7, R \sim 25.8$) imply a $-0.13 \lesssim (V - R)_{0} \lesssim 0.6$. This means that the PSR J1048–5832 spectrum would be unusually red compared to the Vela pulsar. Moreover, the ratio between the unabsorbed optical and X-ray flux of PSR J1048–5832 would be much higher than for other young pulsars. Thus, we conclude that Star D is not the PSR J1048–5832 counterpart. We compared the derived R and V-band upper limits ($R \gtrsim 26.4; V \gtrsim 27.6$) with the extrapolation of the X and γ-ray spectra and constrained the pulsar spectrum at low-energies. In particular, the VLT upper limits suggest that the pulsar spectrum could be consistent with a single power-law, stretching from the γ-rays to the optical.

Key words: Optical: stars – neutron stars

1 INTRODUCTION

PSR J1048–5832 (B1046–58) is a young, Vela-like pulsar in the Carina region, discovered during a 1.4 GHz Parkes survey of the Galactic plane (Johnston et al. 1992). It has a period $P=123.7$ ms and a period derivative $\dot{P} = 9.63 \times 10^{-14}$ s$^{-1}$, which correspond to a characteristic age $\tau_c \sim 20.3$ kyr, a surface dipole magnetic field $B_S \sim 3.5 \times 10^{12}$ G, and a spin down luminosity $E_{SD} \sim 2 \times 10^{36}$ erg s$^{-1}$. The NE2001 model (Cordes & Lazio 2002) yields a distance of 2.7±0.35 kpc, based on the dispersion measure (DM). In the X-rays, PSR J1048–5832 was detected with the Röntgen Satellite (ROSAT) at 0.1–2.4 keV (Becker et al. 1997). Soon after, observations with the Advanced Satellite for Cosmology and Astrophysics (ASCA) found possible evidence of extended X-ray emission (Pivoravoff et al. 2000), possibly associated with a pulsar wind nebula (PWN). High-resolution observations with the *Chandra* X-ray Observatory, indeed revealed the presence of an asymmetric PWN around the pulsar. No X-ray pulsed emission has been detected so far by both *Chandra* and XMM-Newton, and a 3σ upper limit of 53% on the pulsed fraction between 0.5 and 10 keV has been determined (Gonzales et al. 2006). A recent re-analysis of the *Chandra* data of PSR J1048–5832 (Marelli 2012) could separate the pulsar emission from that of its PWN. See Marelli (2012) for details on the observations and data analysis. The combined *Chandra* plus XMM-Newton spectrum of the pulsar is fitted by a power-law (PL) with photon index $\Gamma_X = 1.35 \pm 0.45$ and $N_H = (46.0 \pm 2.3) \times 10^{20}$ cm$^{-2}$, corresponding to an unabsorbed X-ray flux $F_X = 0.490^{+0.183}_{-0.342} \times 10^{-13}$ erg cm$^{-2}$ s$^{-1}$ in the 0.3–10 keV energy range. We note that, using the *Chandra* data, Gonzales et al. (2006) tried to separate the pulsar and PWN components and derived a pulsar photon index $\Gamma_X \sim 2.4$ (quoted without uncertainties) from the fit to the combined *Chandra* plus XMM-Newton spectrum. Our results differ from theirs since we used data from all the XMM-Newton detectors (both the EPIC-PN and MOS) and fitted the combined spectrum using the C-statistic approach in XSPEC. This allowed us to more efficiently separate the pulsar and PWN components in our spectral analysis and derive more accurate values for the pulsar photon index. Moreover, we used slightly different radii to extract the pulsar counts: 20′′ instead of 25′′ for the

* Based on observations made with ESO Telescopes at the La Silla Paranal Observatory under programme ID 384.D-0386(A), 386.D-0585(A).
XMM-Newton data, and 2′′ instead of 1″ for the Chandra data. PSR J1048−5832 was initially proposed by Kaspi et al. (2000) as the counterpart of the unidentified γ-ray source 3EG J1048−5840, detected above 400 MeV by the EGRET instrument aboard the Compton Gamma Ray Observatory (CGRO). More recently, the Large Area Telescope (LAT) aboard the Fermi Gamma-ray Space Telescope detected for the first time γ-ray pulsations above 100 MeV from this pulsar (Abdo et al. 2009), highlighting a clear double−ray source 3EG J1048−5832 and its associated rms in (parentheses).

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2 OBSERVATIONS AND DATA REDUCTION

The R-band images of PSR J1048−5832 were obtained with the VLT Antu telescope at the ESO Paranal Observatory between December 4, 2010 and January 5, 2011 (See Tab. 1) and are available in the ESO archive under programme ID 384.D-0386(A) and 386.D-0585(A). All observations were performed in service mode with the FOcal Reducer/low dispersion Spectrograph (FORS2, Appenzeller et al. 1998). FORS2 is equipped with a red-sensitive MIT detector, a mosaic of two 2k×4k CCDs optimised for wavelengths longer than 6000 Å. The pixel size of FORS2 in standard resolution mode is 0″25 (2 × 2 binning), corresponding to a field of view of 8′3 × 8′3 over the CCD mosaic. Observations were performed in IMAGE mode, with standard low gain, normal readout (200 Kpix/s) and standard-resolution mode. For all the observations, the target was positioned on CHIP1. The R\_SPEC filter (\(\lambda = 6550 \, \AA; \Delta \lambda = 1650 \, \AA\)) was used for all observations. In order to minimise the saturation of bright stars and to allow for cosmic ray removal, a series of 23 short exposures of 599.9 s were obtained, for a total exposure of 13797.7 s (Tab. 1). Exposures were taken in dark time and photometric conditions. The average airmass was below 1.5 and the seeing better than 1″, as determined by the Differential Image Motion Monitor (DIMM).

We used the standard packages in IRAF for the bias subtraction and for flat field correction, selecting the closest-in-time bias and twilight flat field images available on the ESO archive. After these standard corrections, we aligned and stacked the reduced science images using the Swarp tool, included in the Scisoft 7.7 suit\(^3\) applying a 3σ filter to remove the hot/cold pixels and cosmic ray hits from the average image. Since all nights were photometric, we computed the photometry calibration of the average image by using the average of the \(R_{\text{SPEC}}\) night zero points (< 0.05 magnitudes rms) and extinction coefficients computed over the four nights and the colour term for the September 2010–March 2011 semester. All these values are computed by the FORS2 pipeline and are available through the FORS2 data-quality control databases\(^4\). For the average science image we assumed the average airmass computed over all the exposures. We estimated that the effects of the airmass variation over the different exposures (0.11 rms) and the variation of the extinction coefficient (0.012 rms) over the four nights introduce an uncertainty of only 0.05 magnitudes on our photometry.

The astrometric solution of the FORS2 frame is based on the coordinates of the guide star used for the pointing. In order to improve the accuracy of the pulsar position on the frame, we recomputed the astrometric solution using the Two Micron All Sky Survey (2MASS) All-Sky Catalog of Point Sources (Skrutskie et al. 2006). Since the brightest stars in the field are saturated, we selected a subsample of 32 fainter, non-saturated reference stars, by excluding those at clearly bad position, e.g. close to the CCD edges. We then measured the star centroids through Gaussian fitting using the Graphical Astronomy and Image Analysis (GAIA) tool\(^5\). We also used GAIA astrometric fitting routines to compute the pixel-to-sky coordinate transformation, which also account for the CCD distortions. The rms of the astrometric fits was \(\sigma_\alpha = 0′′1\) in the radial direction. To this value we added in quadrature the uncertainty \(\sigma_{\alpha,0} = 0′′1\) of the registration of the FORS2 image on the chosen astrometric reference frame. According to Lattanzi et al. (1997), this is \(\sigma_\alpha = \sqrt{n/N}\sigma_{2\text{MASS}}\), where \(N=32\) is the number of stars used to compute the astrometric solution, \(n=9\) is the number of free parameters in the sky−to−image transformation model, and \(\sigma_{2\text{MASS}}\) is the 1σ mean positional error of 2MASS, which depends on the brightness of the reference stars. The selected 2MASS reference stars have magnitudes 15.5 \(\leq K \leq 13\), which corresponds to \(\sigma_{2\text{MASS}} \leq 0′′2\) (Skrutskie et al. 2006). We also accounted for the 0′015 accuracy of the link of the 2MASS coordinates to the International Celestial Reference Frame (Skrutskie et al. 2006).

Thus, we estimated that the overall 1σ positional uncertainty of our FORS2 astrometry is \(\sim 0′′18\). For consistency, we also recomputed the astrometry calibration of the FORS2 V-band image of Migani et al. (2011) using the 2MASS catalog. The resulting rms of the astrometric fits was \(\sigma_\alpha = 0′′16\), leading to an overall 1σ uncertainty on the V-band image of \(\sim 0′′19\).

3 DATA ANALYSIS AND RESULTS

As done in Migani et al. (2011), also in this work, we assumed as a reference both the Chandra (\(\alpha=10^h 48^m 12^s 64; \delta=-58^\circ 32′ 03′′60\) (MJD= 52859; Gonzalez et al. 2006) and radio interferometry (\(\alpha=10^h 48^m 12^s 604; \delta=-58^\circ 32′ 03′′75\) (MJD=50581; Stappers et al. 1999)) positions that have radial positional uncertainties of 0′′55 and 0′′08, respectively. These are shown in Fig.1 overlaid on the co-added FORS2 R-band image (MJD=55566). Our astrometry is

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1. www.eso.org/archive
2. archive.eso.org/asm/ambient-server
4. www.eso.org/qc
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4 DISCUSSION

Star D (V=26.7±0.2) was ruled out by Mignani et al. (2011) as a possible counterpart to PSR J1048−5832 because of its ≈ 3σ offset from the most recent Chandra position of Gonzalez et al. (2006) and an ≈8% chance coincidence probability. However, since the time span between the Chandra (MJD=52859) and the FORS2 V-band (MJD=55237) observations is ≈6.5 years, the actual pulsar position could have been closer to that of Star D at the latter epoch, owing to its unknown proper motion. For instance, for a pulsar transverse velocity of 1200 km s⁻¹, corresponding to a proper motion of 0°.9 yr⁻¹ at a 2.7 kpc distance, the PSR J1048−5832 position could have been closer to that of Star D by ≈0.7. Such a velocity is about three times as large as the average transverse velocity of radio pulsars (Hobbs et al. 2005). Pulsars with extreme velocities do exist, though, such as the Guitar Nebula pulsar, PSR B2224+65, whose transverse velocity could be as high as 1600 km s⁻¹ (Chatterjee & Cordes 2004). Although the probability that PSR J1048−5832 is moving at such an high velocity and exactly along a position angle of ≈210°, i.e. that of the vector connecting the Chandra position to that of Star D, is small, we cannot rule it out a priori. If Star D was the pulsar counterpart, it should then have the same proper motion as hypothesised for the pulsar. We used the FORS2 V and R-band images, separated by ~ 0.9 years, to measure a possible angular displacement of Star D through relative astrometry. We registered the two images using standard tools in IRAF, with an rms accuracy of ~0.05 pixel per coordinate. Unfortunately, the faintness of Star D prevented us to fits its centroid with an accuracy better than 0.2 pixel per coordinate and we could only set a 1σ upper limit of ~0.4 pixel on its displacement. Accounting for the FORS2 pixel scale (0′′.25), this corresponds to an upper limit of ~0′′.11 yr⁻¹ on the Star D proper motion. If Star D was the pulsar counterpart, this would only allow us to rule out a transverse velocity ≥ 1500 km s⁻¹.

We used the colour information on Star D, now available, to further investigate its possible association with PSR J1048−5832. Star D has an observed (V−R)=0.9±0.3. We checked whether this colour would be compatible with a pulsar spectrum. We corrected the observed fluxes of Star D assuming the reddening towards PSR J1048−5832, computed from the hydrogen column density inferred from the X-ray spectral fits. The NH derived from the best fit to the Chandra plus XMM-Newton X-ray spectrum (Marelli 2012) corresponds to a reddening E(B−V)=0.82±0.04 according to the relation of Predehl & Schmitt (1995). This is much lower than the value inferred from the XMM-Newton spectrum alone (Marelli et al. 2011), E(B−V)=1.6±0.7, and assumed by Mignani et al. (2011) in their analysis of the VLT data. Using the interstellar extinction coefficients of Fitzpatrick (1999), we derived AV=2.55±0.13 and AR=1.91±0.09. Thus, if Star D were affected by the same extinction as the pulsar, its intrinsic colour would be (V−R)₀=0.26±0.33, i.e. its spectrum would be unconstrained. We note that the Vela pulsar has (V−R)₀≈0.3 (Mignani & Caraveo 2001), where the uncertainty due to the reddening correction for NH = (2.2±0.5)×10²⁰ cm⁻² (Marelli et al. 2011) is marginal, i.e. bluer than it would be for PSR J1048−5832. However, Vela has a flat PL spectrum in the optical (Mignani et al. 2007), whereas some pulsars might have steeper spectra that rise towards longer wavelengths, such as PSR B0540−69 (Mignani et al. 2010a), which would yield colours redder than Vela. Thus, the colour of Star D might not be incompatible with that of a pulsar.

We note that the flux of Star D would be incompatible with the expected optical emission of PSR J1048−5832 extrapolated from the luminosity of the Vela pulsar. PSR J1048−5832 has a spin-down luminosity ESD ≈ 2 × 10³⁶ erg s⁻¹, a factor of 3 smaller than the Vela pulsar (ESD ≈ 6.9 × 10³⁶ erg s⁻¹). Moreover, it is at a larger distance (2.7±0.35 kpc) than Vela (0.287±0.02 kpc; Dodson et al. 2003) and is affected by a factor of ~20 larger reddening. Assuming that the PSR J1048−5832 optical luminosity Lopt is the same fraction of the spin-down luminosity ESD as the Vela pulsar and accounting for the distance and reddening would then yield V ≈ 32.2. If Star D were the PSR J1048−5832 counterpart, this would emit in the optical up to ≈ 3.7 × 10⁻⁷ of its spin-down luminosity. This value is ≈100 larger than the Vela pulsar, and comparable to that of middle-aged and older pulsars (Zharkov et al. 2006). However, the evolution of the pulsar optical emission efficiency, ηopt ≡ Lopt/ESD, with the characteristic age is quite

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Figure 1. VLT/FORS2 image of the PSR J1048−5832 field (R band; 13797.7 s). The Chandra (Gonzalez et al. 2006) and radio interferometry (Stappers et al. 1999) positions are shown as circles. The star closest to the Chandra position is Star D of Mignani et al. (2011).

http://archive.eso.org/wdb/wdb/eso/abstract/query?&prodigid=386.D-0585(A)
 uncertain. In particular, this is true for spin-down ages between 10 and 100 krs, where Vela is the only pulsar identified in the optical (Mignani 2011). Thus, we do not know whether the sharp decrease of $\eta_{\text{opt}}$ observed for Vela, which is only $\approx 5$ times older than the Crab but has a $\approx 1000$ times lower efficiency, is peculiar of this object or is representative of all Vela-like pulsars. Indeed, the upper limits on the optical luminosity inferred for other Vela-like pulsars (e.g., Mignani et al. 1999; 2011; 2012) imply emission efficiencies $\eta_{\text{opt}} \lesssim 10^{-6}$, which would be consistent with that derived for PSR J1048—5832 if Star D was its optical counterpart.

We investigated this association from the corresponding X-ray-to-optical flux ratio $F_X/F_V$, where $F_X$ and $F_V$ are the unabsorbed Chandra and VLT fluxes of PSR J1048—5832 and Star D, respectively. For the assumed reddening, the association with Star D would imply, at most, an unabsorbed optical flux $F_V \sim 8.6 \times 10^{-16}$ erg cm$^{-2}$ s$^{-1}$ and $F_X/F_V \sim 100$. This value is much lower than expected from the comparison with other young rotation-powered pulsars (e.g., Zharikov et al. 2006), for which the non-thermal, unabsorbed X-ray-to-optical flux ratio is always larger than $\approx 800$. Thus, we conclude that Star D cannot be the PSR J1048—5832 counterpart. The R-band upper limit implies a flux $F_X^R \lesssim 0.46 \mu$Jy, corrected for the interstellar extinction. This flux value is quite close to that derived in the V band ($F_X^V \lesssim 0.34 \mu$Jy). The optical flux upper limit on PSR J1048—5832 imply, for the revised value of the interstellar absorption, an optical emission efficiency $\eta_{\text{opt}} \lesssim 1.9 \times 10^{-7}$, accounting for the uncertainty on the $N_I$ and the pulsar distance.

We compared the optical flux upper limits with the low-energy extrapolation of the X-ray and $\gamma$-ray spectra (Fig. 2). For the X-rays, we assumed the spectral model that fits the Chandra plus XMM-Newton data (Marelli 2012), a PL with photon index $\Gamma_X=1.35\pm0.45$, while for the $\gamma$-rays we assumed the PL with photon index $\Gamma_{\gamma}=1.38\pm0.13$ and exponential cut-off at $\sim 2.3$ GeV that fits the Fermi/LAT data (Abdo et al. 2009). As seen from Fig. 2, the extrapolation of the $\gamma$-ray PL can be consistent with the X-ray PL spectrum. This changes the conclusions reported in Mignani et al. (2011), based on the XMM-Newton spectrum, where the presence of a spectral break between the X-rays and the $\gamma$-rays was apparent. Fig. 2 also shows that the optical flux upper limits are just above the extrapolation of the X-ray and $\gamma$-ray spectra. This means that the multi-wavelength spectrum of PSR J1048—5832 could be consistent with a single PL stretching from the $\gamma$-rays to the optical (see also Durant et al. 2011), at variance with other Fermi pulsars (e.g., Mignani et al. 2010b; 2011; 2012), where there is evidence for spectral breaks between the PL components in the different energy bands, which are possibly representative of different energy and particle distributions. Thus, the detection of PSR J1048—5832 in the optical would be crucial to confirm this scenario.

5 CONCLUSIONS

The new VLT observations of PSR J1048—5832 confirm that the object detected closest to the Chandra position (Star D) is an unrelated field star. No other possible counterpart is detected down to $R \sim 26.4$ and $V \sim 27.6$. Thus, PSR J1048—5832 remains unidentified in the optical. It is clear that its detection is close to the sensitivity limit of 8m-class telescopes. Observations at longer wavelength might be more successful, though, were the optical/infrared spectrum of PSR J1048—5832 follow the extrapolation of the X-ray PL. An updated Chandra or radio position for PSR J1048—5832 would allow one to minimise the position uncertainty due to the unknown pulsar proper motion and would represent an advantage in the search for its optical/infrared counterpart.

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