



Astronomical calibration of the Ypresian timescale: implications for seafloor spreading rates and the chaotic behavior of the solar system?

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Abstract. To fully understand the global climate dynamics of the warm early Eocene with its reoccurring hyperthermal events, an accurate high-fidelity age model is required. The Ypresian stage (56–47.8 Ma) covers a key interval within the Eocene as it ranges from the warmest marine temperatures in the early Eocene to the long-term cooling trends in the middle Eocene. Despite the recent development of detailed marine isotope records spanning portions of the Ypresian stage, key records to establish a complete astronomically calibrated age model for the Ypresian are still missing. Here we present new high-resolution X-ray fluorescence (XRF) core scanning iron intensity, bulk stable isotope, calcareous nannofossil, and magnetostratigraphic data generated on core material from ODP Sites 1258 (Leg 207, Demerara Rise), 1262, 1263, 1265, and 1267 (Leg 208, Walvis Ridge) recovered in the equatorial and South Atlantic Ocean. By combining new data with published records, a 405 kyr eccentricity cyclostratigraphic framework was established, revealing a 300–400 kyr long condensed interval for magnetochron C22n in the Leg 208 succession. Because the amplitudes are dominated by eccentricity, the XRF data help to identify the most suitable orbital solution for astronomical tuning of the Ypresian. Our new records fit best with the La2010b numerical

solution for eccentricity, which was used as a target curve for compiling the Ypresian astronomical timescale (YATS). The consistent positions of the very long eccentricity minima in the geological data and the La2010b solution suggest that the macroscopic feature displaying the chaotic diffusion of the planetary orbits, the transition from libration to circulation in the combination of angles in the precession motion of the orbits of Earth and Mars, occurred ~ 52 Ma. This adds to the geological evidence for the chaotic behavior of the solar system. Additionally, the new astrochronology and revised magnetostratigraphy provide robust ages and durations for Chrons C21n to C24n (47–54 Ma), revealing a major change in spreading rates in the interval from 51.0 to 52.5 Ma. This major change in spreading rates is synchronous with a global reorganization of the plate–mantle system and the chaotic diffusion of the planetary orbits. The newly provided YATS also includes new absolute ages for biostratigraphic events, magnetic polarity reversals, and early Eocene hyperthermal events. Our new bio- and magnetostratigraphically calibrated stable isotope compilation may act as a reference for further paleoclimate studies of the Ypresian, which is of special interest because of the outgoing warming and increasingly cooling phase. Finally, our approach of integrating the

complex comprehensive data sets unearths some challenges and uncertainties but also validates the high potential of chemostratigraphy, magnetostratigraphy, and biostratigraphy in unprecedented detail being most significant for an accurate chronostratigraphy.

1 Introduction

The Ypresian stage, from 56.0 to 47.8 Ma, represents the first ~ 8 million years of the Eocene epoch (Vandenberghé et al., 2012), an interval characterized by the warmest deep-sea temperatures of the Cenozoic era (Zachos et al., 2008), multiple transient global warming events (Cramer et al., 2003; Lourens et al., 2005; Agnini et al., 2009; Galeotti et al., 2010; Leon-Rodríguez and Dickens 2010; Stap et al., 2010; Zachos et al., 2010; Sexton et al., 2011; Slotnick et al., 2012; Littler et al., 2014; Kirtland Turner et al., 2014; Lauretano et al., 2015, 2016), and major faunal turnovers (Thomas and Shackleton, 1996; Gingerich, 2003; Clyde et al., 2007). Climatic records from the Ypresian greenhouse are of special interest because they potentially capture the behavior of Earth's climate system under $p\text{CO}_2$ concentrations likely to be reached in the near future (Zachos et al., 2008; Meinshausen et al., 2011). For the detailed reconstruction of the Eocene greenhouse climate system, a complete, precise, and highly accurate stratigraphic framework to determine rates of climatic processes and timing of events is required.

Despite the recent progress in construction of astronomically calibrated age models for the Eocene using geochemical records (Cramer et al., 2003; Lourens et al., 2005; Westerhold et al., 2007; Galeotti et al., 2010; Jovane et al., 2010; Westerhold et al., 2012, 2014, 2015; Westerhold and Röhl, 2013; Littler et al., 2014; Lauretano et al., 2016), a fully consistent astrochronology for the Ypresian is not yet available. Two major issues have to be solved to achieve a complete Ypresian astronomical timescale (YATS): (1) the “50 Ma discrepancy” in magnetostratigraphy (Vandenberghé et al., 2012; Westerhold et al., 2015) and (2) the exact number of 405 kyr eccentricity cycles in magnetochron C23 (Lauretano et al., 2016).

The 50 Ma discrepancy arises from the short duration of magnetochron C23n.2n in Ocean Drilling Program (ODP) Site 1258 data (Suganuma and Ogg, 2006; Westerhold and Röhl, 2009) that results in a very unlikely abrupt global increase in spreading rate for this chron only at around 50 Ma (Vandenberghé et al., 2012). Assuming a 400 kyr longer duration for Chron C22r in the same record would result in lower spreading rates than calculated on the basis of the standard geomagnetic polarity timescale (GPTS) using the synthetic magnetic anomaly profile of the South Atlantic (Cande and Kent, 1992, 1995; CK95). However, the Site 1258 magnetostratigraphy is currently the only available record covering the entire Ypresian. Paleomagnetic interpretation for sediments from Site 1258 is difficult because the core was

retrieved by rotary drilling. Minicores from Site 1258 gave a relatively good magnetostratigraphy but it is based on a rather subjective method of polarity interpretation (Suganuma and Ogg, 2006). After revision of the composite record of Site 1258, the interpretation was improved (Westerhold and Röhl, 2009), but the interpretation for Chrons C22r, C23n, and C23r remained ambiguous.

Eccentricity-modulated precession cycles in X-ray fluorescence (XRF) core scanning iron (Fe) intensities from Site 1258 and Leg 208 sites helped to set up the first cyclostratigraphic age models for the Ypresian (Westerhold et al., 2007; Westerhold and Röhl 2009). High-resolution bulk and benthic stable isotope data for the early (Cramer et al., 2003; Zachos et al., 2010; Stap et al., 2010; Littler et al., 2014; Lauretano et al., 2015, 2016), middle (Kirtland Turner et al., 2014), and late Ypresian (Sexton et al., 2011) showed exceptionally strong 100 and 405 kyr eccentricity cycle variations that were partly utilized for age model construction. In order to test the 50 Ma discrepancy and the astronomical age model at ODP Site 1258, high-resolution benthic stable isotope records spanning 54 to 49 Ma were compiled from ODP Site 1262 (Littler et al., 2014; Lauretano et al., 2015) and Site 1263 (Stap et al., 2010; Lauretano et al., 2015, 2016) and were astronomically tuned to the La2010d (Laskar et al., 2011a) orbital solution (Lauretano et al., 2016). In theory, the dominant eccentricity-related cyclicity in this interval should enable development of a robust astrochronology. However, a period of low-amplitude variability and a major shift in benthic $\delta^{13}\text{C}$ data in Chron C23n and upper Chron C23r complicate the identification of the correct number of 405 kyr cycles. The favored age model in Lauretano et al. (2016) proposes two 405 kyr cycles in this complex interval, which provided absolute age estimates for early Eocene hyperthermal events. The two 405 kyr cycle age model is favored because it assumes constant sedimentation rates. However, it did not solve the 50 Ma discrepancy because the duration for C23n.2n is much too short (295 kyr) compared to CK95 (696 kyr) in ODP 1258.

To establish a complete YATS with consistent GPTS ages, a new deep-sea magneto-cyclostratigraphic record is needed to test the ODP Site 1258 magnetostratigraphy. In particular, the durations of Chrons C22 and C23 have to be evaluated. Here we present a new complete magnetostratigraphy spanning Chron C21n to C24n by integrating records from ODP Leg 208 Sites 1262, 1263, 1265, and 1267. New XRF core scanning data and core images are used for ultra-high-resolution correlation between the Leg 208 sites. New calcareous nannofossil data from Sites 1258, 1263, and 1265 are presented and combined with published datums for direct calibration to magnetostratigraphy and revision of age datums. After integration of Leg 208 and ODP Leg 207 Site 1258 isotope data with new XRF data and core images, a cyclostratigraphic framework was compiled that was subsequently astronomically tuned to the La2010 (Laskar et al., 2011a) orbital solution. This study provides the first com-

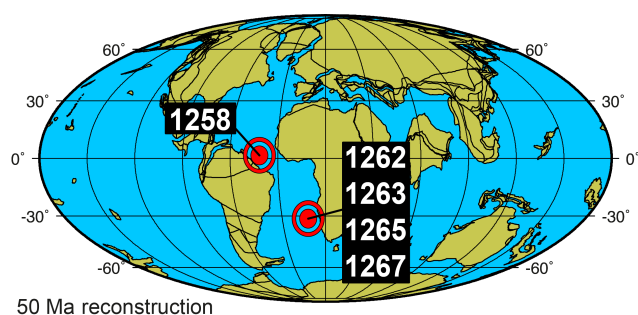


Figure 1. Location of ODP Site 1258 and Leg 208 sites on a 50 Ma paleogeographic reconstruction in Mollweide projection (from <http://www.ods.n.de>).

plete and consistent bio-, chemo-, and magnetostratigraphy for the Ypresian stage. Our new data sets combining two regional records of ODP Legs 207 and 208 reveal unprecedented insight into possible challenges and uncertainties but also demonstrate the high potential of detailed chemostratigraphy, magnetostratigraphy, and biostratigraphy that will be the prerequisite for the major field of chronostratigraphy.

2 Material and methods

New data for this study were generated from sediment cores retrieved during ODP Leg 207 (Site 1258; Erbacher et al., 2004) and Leg 208 (Sites 1262, 1263, 1265, 1267; Zachos et al., 2004). The carbonate-rich sequences drilled at the equatorial Atlantic Demerara Rise (Leg 207) and the southeast Atlantic Walvis Ridge (Leg 208, Fig. 1) were recovered from multiple holes combined with composite records. For Sites 1262 and 1265 the shipboard composite depth was used (mcd – meters composite depth) for analysis (Zachos et al., 2004). For Sites 1258, 1263, and 1267 the revised composite record was applied (rmcd – revised composite depth; Westerhold et al., 2007, 2015; Westerhold and Röhl, 2009) for further analysis.

2.1 XRF core scanner data

To obtain a complete record for the Ypresian, published data had to be combined with the new XRF Fe core scanning data. The newly acquired data presented here have been measured on the three XRF core scanners at MARUM – Center for Marine Environmental Sciences, University of Bremen, with various hardware generations and under different settings (for details see Supplement). Site 1262 XRF data were obtained from 92.74–104.57 and 111.10–112.40 mcd, and were combined with data from Westerhold et al. (2007, 2012). Site 1263 data reported here are from 168.08 to 289.00 rmcd and were combined with data from Westerhold et al. (2007). Site 1265 XRF data were measured from 228.10 to 275.00 mcd and merged with data from Wester-

hold et al. (2007). New Site 1267 XRF data were generated from 153.12 to 236.53 rmcd. Settings for XRF core scanning of Site 1258 are given in Westerhold and Röhl (2009). Here we report the data from 134.61 to 212.49 rmcd for Site 1258. Combined with published data, 18 000 new XRF core scanning data points (see Supplement) result in a total of more than 30 000 data points covering the latest Paleocene and Ypresian. This enormous data set is the prerequisite for correlating the five drill sites in detail.

2.2 Bulk stable isotope data

To close a stratigraphic gap in the stable isotope record for Site 1263, a section (230.50 to 239.00 rmcd) of nearby Site 1265 was selected. A total of 369 sediment samples were collected, freeze-dried, and powdered (Table S8 in the Supplement). The $\delta^{13}\text{C}$ and $\delta^{18}\text{O}$ values were measured at MARUM – Center for Marine Environmental Sciences, University of Bremen, and are reported relative to the Vienna Pee Dee Belemnite international standard, determined via adjustment to calibrated in-house standards and NBS-19. Analyses at MARUM were carried out on a Finnigan MAT 251 mass spectrometer equipped with an automated carbonate preparation line (Kiel III). The carbonate was reacted with orthophosphoric acid at 75 °C. Analytical precision based on replicate analyses of an in-house standard (Solnhofen Limestone) averaged 0.03 ‰ (1σ) for $\delta^{13}\text{C}$ and 0.06 ‰ (1σ) for $\delta^{18}\text{O}$ for samples measured in 2014 and 0.04 ‰ for $\delta^{13}\text{C}$ and 0.04 ‰ for $\delta^{18}\text{O}$ for samples run in 2015.

Published stable isotope data also used in this study have been compiled from Lourens et al. (2005), Zachos et al. (2005, 2010), McCarren et al. (2008), Stap et al. (2009, 2010), Sexton et al. (2011), Westerhold et al. (2012, 2015), Kirtland Turner et al. (2014), Littler et al. (2014), and Laurentano et al. (2015, 2016). All data are provided in the Supplement (and PANGAEA) relative to the Site 1263 depth and with the tuned age added.

2.3 Paleomagnetic data Leg 208

Natural remanent magnetization (NRM) was measured on 400 discrete cube samples (dimensions 2 cm × 2 cm × 2 cm) to document magnetic polarity boundaries C20r to C24r at ODP Sites 1262 (64), 1263 (128), 1265 (115), and 1267 (89). Discrete samples were analyzed at the Faculty of Geosciences, University of Bremen. Paleomagnetic directions and magnetization intensities were measured on a cryogenic magnetometer (2G Enterprises model 755 HR). NRM was measured on each sample before being subjected to a systematic alternating field demagnetization treatment involving steps of 7.5, 10, 15, 20, 25, 30, 40, and 60 mT. Intensities of orthogonal magnetic components of the remanent magnetization were measured after each step. Raw inclination, declination, and intensity data for each measurement step are provided in Tables S17–S20 in the Supplement. The calcu-

lated characteristic remanent magnetization in Tables S21–S25 and the magnetostratigraphic interpretations, including published data, are recorded in Tables S26–S30 for each of the five sites.

2.4 Calcareous nannofossils at ODP Sites 1258, 1263, and 1265

Smear slides were processed following the standard procedures described in Bown et al. (1998) in order to investigate calcareous nannofossil assemblages. High-resolution semiquantitative counting methods, which consist of counting the number of forms ascribed to the same taxon detected in or normalized to a prefixed area (i.e., 1 mm²; Backman and Shackleton, 1983; Agnini et al., 2014), were used. These counts provide very detailed abundance patterns, which permit the precise identification of the position of each biohorizon, especially in those cases in which either the marker species displays an overall uneven distribution or a marker species is rare and discontinuous at the base and top of its stratigraphic range. On this basis and according to Agnini et al. (2014), six types of biohorizons are adopted, these are base rare (Br), base (B), base common and continuous (Bc), top (T), top common and continuous (Tc), and crossover (X). In two cases we used the increase–decrease in abundance of taxa (i.e., *Fasciculithus* spp., *Zygrhablithus bijugatus*) to define additional biohorizons. Taxonomic concepts adopted in this study are those of Perch-Nielsen (1985) and Agnini et al. (2007).

Calcareous nannofossil biostratigraphic analyses were newly performed or refined for the early Eocene of ODP Sites 1258, 1263, and 1265. The biohorizons used in different low-latitude to midlatitude zonations (Martini, 1971; Okada and Bukry, 1980; Agnini et al., 2014), as well as additional biohorizons, were recognized. A set of 168 samples were analyzed from the Demerara Rise (equatorial Atlantic) which allows for the identification of 40 biohorizons. At Walvis Ridge (SE Atlantic), a total of 181 samples were studied from ODP Sites 1263 (77) and 1265 (104), which permits the detection of 27 biohorizons. New data were integrated with ship data or other published results to obtain a more complete and reliable data set. Tables of the calcareous nannofossil biohorizons are given in Tables S31–S35.

3 Results

All data generated within and available data compiled for this study are combined in the data set file and available open access online at <http://doi.pangaea.de/10.1594/PANGAEA.871246>.

3.1 XRF core scanning results

Fe intensity data reveal the cyclic pattern commonly observed for the interval from PETM to ETM-2 (Westerhold

et al., 2007; Littler et al., 2014), with higher values in darker, more clay-rich layers. A decrease in carbonate content around 51 Ma for all sites (Zachos et al., 2004) is reflected by overall higher Fe intensities. Multiple distinct peaks in Fe data for Chrons C22 and C23n correspond to strikingly bundled sets of clay-rich dark intervals for all sites. At Site 1262 the XRF Fe record ends with the shoaling of the carbonate compensation depth (CCD) above the site in Chron C21r around 93 mcd (Zachos et al., 2004). Generally, the records from Sites 1262 and 1267 are of lower resolution than those from Sites 1263 and 1265 due to the regional decline in carbonate accumulation rates with increasing water depth. Site 1263 shows the most persistent high-resolution XRF Fe intensity signal. A gap in the 1263 record from 229.15 to 233.68 mcd caused by a drilling disturbance (Zachos et al., 2004) can successfully be bridged by the records from Sites 1265 and 1267. The high-resolution XRF Fe data show consistent patterns that are required to do a detailed site-to-site correlation and integration of Leg 208 Sites 1262, 1263, 1265, and 1267. New XRF Fe intensity data for Site 1258 between PETM and ETM-2 reveal the same eccentricity-modulated precession cycles as observed from Leg 208 sites.

XRF core scanning Fe intensity data from four Leg 208 sites are shown in Fig. S1 in the Supplement and given in Tables S1 to S7. All data plotted versus Site 1263 depth are given in Fig. 2 from ETM-2 to Chron C20r. Data for the interval from the PETM to ETM-2 are plotted in Fig. S7 in the Supplement. Due to the large and very detailed data set and the fact that most of the data from PETM to ETM-2 have been published previously, priority for figures in the main paper is on the interval from ETM-2 to Chron C20r.

3.2 Bulk stable isotope results

To obtain a complete bulk stable isotope record for Leg 208 sites, the gap in the Site 1263 record was bridged by incorporating bulk data from Site 1265 (Fig. S1a). Bulk stable isotope data from Site 1265 show cyclic variations between 1.6 and 2.2 ‰ and match with overlapping data from Site 1263 (Fig. 2; Westerhold et al., 2015, 217 to 227 mcd of 1263). As observed previously (Zachos et al., 2010; Littler et al., 2014), lighter bulk $\delta^{13}\text{C}$ data coincide with dark clay-rich intervals with relatively higher XRF Fe intensities at all Leg 208 sites. All bulk and benthic data compiled for this study are presented in Tables S8 to S16.

3.3 Core image processing and site-to-site correlation

To correlate and integrate Leg 208 and Site 1258 records, the new software tool CODD (Code for Ocean Drilling Data; Wilkens et al., 2017) was utilized. This tool greatly facilitates handling of large and complex data sets and allows the use of core images for scientific analysis. For all sites in the study, core images and all available data were assembled by

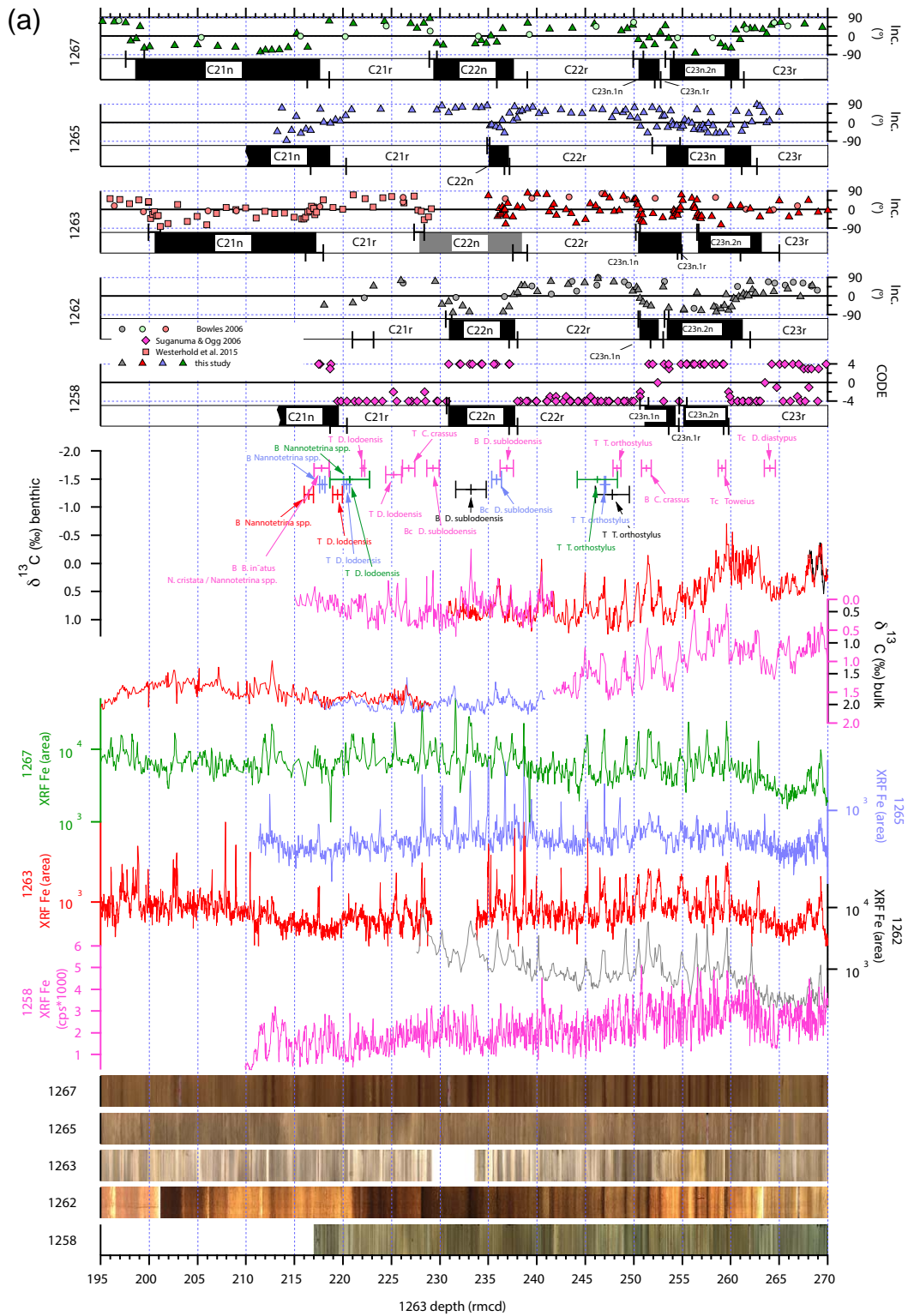


Figure 2.

holes. Then the composite records were cross-checked and assembled. In order to be able to use data generated outside of the splice, all cores were mapped onto the splice using differential stretching and squeezing (Tables S36–S40). Site 1263 was chosen as the reference site to correlate all other sites to because it has the highest sedimentation rates and the most detailed stable isotope data (Lauretano et al., 2015, 2016; McCarren et al., 2008; Stap et al., 2010; Westerhold et al., 2015). XRF Fe and core images primarily guided the correlation between sites (Fig. 2). Stable isotope data are used to assess the correlation between Sites 1258 and 1263 because the XRF Fe data of Site 1258 are dominated by precession cycles and thus difficult to directly correlate to the eccentricity-dominated cycles at Site 1263 (Westerhold and Röhl, 2009). Figure 2 shows the outstanding match between the sites just by visual comparison of the core images. Existing correlations between Leg 208 sites (Röhl et al., 2007; Westerhold et al., 2007; Lauretano et al., 2015) were updated and adjusted if needed. Correlation tie points are provided in Tables S41–S44. The recent correlation between Sites 1262 and 1263 as well as Sites 1263 and 1258 by Lauretano et al. (2015, 2016) based on stable isotope data were further refined as well. The primary modifications between Sites 1262 and 1263 were made in a short interval (265 to 267 rmc) of Site 1263. The detailed comparison with Site 1265 reveals a gap of about three precession cycles in XRF Fe at 284.40 rmc at Site 1263 due to a core break (Fig. S6). No additional mismatches were recognized, suggesting that the combined Leg 208 records represent the complete stratigraphic sequence for Walvis Ridge. However, correlation of 1258 to 1263 shows overall good agreement except for the interval from 230 to 235 rmc of 1263. Fine-scale comparison of the benthic Site 1258 and bulk Site 1265 stable $\delta^{13}\text{C}$ data show that Leg 208 sediments encompass a regionally condensed interval at $\sim 229\text{--}230$ rmc of Site 1263 (Fig. S8). The missing stratigraphic interval spans 300 to 400 kyr as depicted from the Site 1258 benthic record and thus needs to be corrected for in the Site 1263 astronomically tuned age model. This demonstrates the benefits of utilizing multiple records from different regions with robust composite records to establish a highly accurate stratigraphic framework based on orbital tuning for any given interval.

3.4 Magnetostratigraphic results and interpretation

Vector analysis according to the method by Kirschvink (1980) without anchoring to the origin of the orthogonal projections was applied to the results of the alternating field demagnetization of NRM to determine the characteristic remanent magnetization (ChRM). The maximum angular deviation (MAD) values were computed, reflecting the quality of individual magnetic component directions. Most of the MAD values are below 10 (Fig. S4, Tables S17–S30). Figure S5 displays the demagnetization characteristics of a sample with reversed polarity from

C22r and a sample with normal polarity from C22n. As an example of samples with demagnetization behavior with larger scatter (larger MAD), data from a sample at the C22n–C22r reversal are plotted in Fig. S5. The larger MADs that can be identified at Leg 208 sites in a few samples are not simply related to the intensity of their remanent magnetization. The median destructive field (MDF) of the NRM demagnetization is comparably low for most of the samples. It ranges from 2.6 to 24 mT (mean 6.1 ± 3.8 mT), indicating a magnetically soft overprint in many samples. The interpretation of the ChRM in terms of magnetic polarity is focused on the inclination data, which provide a reliable magnetostratigraphy for most intervals.

Recognition of calcareous nannofossil events allows for the magnetic chrons to be clearly identified from C20r to C24r (Figs. 2 and S4, Tables S26–S30). Raw inclination, declination, and intensity data for each measurement step for Leg 208 sites are given in Tables S17–S20. Magnetostratigraphic interpretation is given in Tables S21–S25. Processed paleomagnetic data from Leg 208 sites, the basis for the magnetostratigraphic interpretation, are provided in Tables S26–S30 for each site. The assignment of error bars, as with all magnetostratigraphic data, is a subjective endeavor. The error bar for Leg 208 data marks the interval in which the inclination shifts from clearly reversed to clearly normal polarity or vice versa. Poorer sample resolution and/or ambiguous or transitional inclination values across a reversal will thus increase the error bar. We did not apply an inclination threshold value to mark a shift in polarity because the reversals occur at different seafloor depths at all sites. Drilling depth and compaction difference between sites might have affected the inclination at each site differently. A much more sharply defined length of error bars could be derived from higher-resolution data (e.g., by analyzing u channels), which is beyond the scope of this study.

Having multiple magnetostratigraphic records from the same region combined with the established high-resolution correlation allows the quality of the paleomagnetic data to be evaluated and inconsistencies to be identified. This, again, is crucial for resolving the 50 Ma discrepancy because a single magnetostratigraphic record from one succession could contain significant unresolved errors. Plotting all Leg 208 ChRM data and the published Site 1258 magnetostratigraphic interpretation against Site 1263 depth immediately shows how consistent but also dynamic the magnetostratigraphy of each site can be (Fig. 2). For example, Chron C22n is clearly too short at Site 1265, which could be related to the condensed interval in this part. Sites 1262 and 1267, however, are consistent with the Site 1258 Chron C22n thickness. At Site 1263 the top of Chron C22n is compromised by drilling disturbance, and the base of Chron C22n is spread over a larger interval than at the other Leg 208 sites. Chron C21n is well captured at Sites 1258, 1263, and 1267, and the base at Site 1265 is also captured well. The top of Chron C23n is consistent between Sites 1258, 1262, 1263, and 1267. The signal

from Site 1265 is a bit noisy and a clear identification for the top of Chron C23n is difficult. The normal interval labeled as C23n could probably only be C23n.2n. The bottom of Chron C23n is consistent within error at Sites 1262, 1265, and 1267, with Site 1262 giving the best signal. The ChRM of Site 1263 does not provide an interpretable signal below 260 rmd, preventing the identification of the base of C23n and the entire Chron C24n. Clearly, comparison to Site 1258 reveals that Chron C23n is too short at Site 1258, probably due to the position of the base Chron C23n. Chron C24n can be identified in Leg 208 but has relatively larger error bars than the other chrons. The top of Chron C24n spreads out from ~270 to 273 rmd of Site 1263 considering all sites. Taking the overlaps of the error bars into account, the best position for the top and bottom of Chron C24n is taken from Site 1262. More difficult to determine is the exact position of the reversals in Chrons C24n, C24n.1r, and C24n.2r. No data are available for Sites 1263 and 1265 in this interval. Resolution at Site 1262 is too low to resolve the short reversed chrons. Sites 1267 and 1258 do not give consistent results either. For the moment, the Site 1258 positions are used for the combined magnetostratigraphy but will need future revision. Based on the integration of all data and evaluation of errors, a magnetostratigraphy for the Ypresian was constructed and is given in Table S45. The results of the combined ChRM data on the high-resolution correlation suggest that a magnetostratigraphic interpretation from a single site might contain significant errors that need to be taken into account. Thus, a magnetostratigraphic interpretation from a single site or location can lead to major discrepancies when used as a template for orbital tuning. The new multi-site data already resolve the 50 Ma discrepancy showing that Chron C23n is too short at Site 1258, causing spreading rates that are too high for the South Atlantic. All of the uncertainties above have been considered while doing the time series analysis and subsequent astronomical calibration for the Ypresian.

3.5 Calcareous nannofossil events at Sites 1258, 1263, and 1265

High-resolution correlation between the sites allows us to investigate how reliably calcareous nannofossil datums can be determined, especially over the depth transect of Leg 208. Therefore, key taxa were targeted to be identified at Sites 1263 and 1265 to be compared to the high-resolution work at Site 1262 (Agnini et al., 2007) and low-resolution shipboard data at Site 1267. Biostratigraphic datums are transformed in biochronological data using the integrated bio-, magneto-, and astrocyclostratigraphic age model developed in this study (Tables S31–S35). Age estimations of calcareous nannofossil biohorizons are generally consistent through the Walvis Ridge sites and in agreement with recently published biochronological data (see Agnini et al., 2014, for review). Most of these biohorizons, in particular almost all the bioevents used in previous and more recent calcareous nan-

nofossil biozonations, were proven to represent reliable data and powerful tools for highly resolved correlations. A total of 18 biohorizons across the study interval (i.e., decrease in diversity of *Fasciculithus* spp., B of *Rhomboaster* spp., the crossover (X) between *Fasciculithus* spp. and *Zyghrablithus bijugatus*, the T of *Fasciculithus* spp., the base of *D. diastypus*, the T of *Tribrachiatus orthostylus*, the T of *Tribrachiatus contortus*, the Tc of *Discoaster multiradiatus*, the B of *Sphenolithus radians*, the B of *Girgisia gammation*, the T of *D. multiradiatus*, the Br and B of *Discoaster lodoensis*, the B of *Chiphragmalithus* spp. (circular), the T of *T. orthostylus*, the T of *D. lodoensis*, the B of *Chiphragmalithus calathus*, and the B of *Nannotetrina* spp.) were calibrated, and the age estimations are impressively consistent throughout Walvis Ridge sites. Two exceptions are the B of *Discoaster sublodoensis*, for which the estimates calibrated at different sites show a high degree of uncertainty, and the B of *Blackites inflatus*, which was not possible to identify because of the absence of this taxon at Walvis Ridge. Other biohorizons, as for instance the B of *Coccolithus crassus*, are not tested at Walvis Ridge but seem to be promising if data from Demerara Rise and data from the Tethyan realm are compared (Agnini et al., 2014).

Data from ODP Site 1258 were produced to biostratigraphically frame the study succession but these data are also used to investigate the degree of reliability of calcareous nannofossil data over wide areas. In general, the stratigraphic positions as well as the ranking and spacing of the biohorizons detected at this site are in fair agreement with data from Walvis Ridge. The two biochronological data sets presented for Walvis Ridge and Demerara Rise sites showed that the ages calculated for some biohorizons, the B of *Rhomboaster* spp., the T of *Fasciculithus* spp., the B of *T. orthostylus*, the Tc and T of *D. multiradiatus*, the B of *D. lodoensis*, the T of *T. tribrachiatus*, and the B of *Nannotetrina* spp., are in fact quite close. However, some other results (i.e., the B of *D. diastypus* and the B of *T. contortus*) are certainly anticipated with respect to the results obtained at Walvis Ridge. These discrepancies are of particular relevance in the mid-upper part of Chron C24r, though a general slight offset is observable between the two areas. In accordance with the age models developed for the study sites, the anticipation of the first occurrences of some taxa could be explained as related to the warm water preferences of the taxa considered (e.g., *Discoaster* and *Tribrachiatus*), but more data are needed to confirm if the diachroneity recorded at Demerara Rise is a general feature of the equatorial latitudes or rather, and more likely, something controlled by local conditions.

4 Astronomical calibration of the Ypresian

Time series analysis of early Eocene records that are used here already showed that the dominant cyclicality in multiple proxy records is related to eccentricity. The interval from

PETM to ETM-2 is dominated by eccentricity-modulated precession cycles that are an impressive recorder of Earth orbital variations through time and the climatic response to them (Lourens et al., 2005; Zachos et al., 2010; Littler et al., 2014). Data from this interval not only allowed the construction of high-precision cyclostratigraphies (Lourens et al., 2005; Westerhold et al., 2007) but also the test of these astrochronologies as well as the theoretical astronomical solutions (Westerhold et al., 2007, 2012; Meyers, 2015). As observed in a compilation of late Paleocene to early Eocene stable isotope data, the $\delta^{13}\text{C}$ variations in both bulk and benthic records from this time show the clear imprint of eccentricity variations, with lighter values occurring in eccentricity maxima (Cramer et al., 2003; Lourens et al., 2005; Zachos et al., 2010; Littler et al., 2014). Hence, bulk and benthic stable isotope data helped to develop astrochronologies from the PETM up to ~ 49 Ma spanning Chrons C22r to C24r at Leg 208 sites (Zachos et al., 2010; Littler et al., 2014; Lauretano et al., 2015, 2016). For Site 1258 a first cyclostratigraphic age model based on XRF core scanning Fe intensity was made from ETM-2 to the base of Chron C21n (~ 47 Ma) leading to revised estimates for the reversal ages from Chrons C24n to C21n (Westerhold and Röhl, 2009). Because of higher sedimentation rates than observed at Leg 208 sites, cyclicity in the Site 1258 XRF Fe data is mainly precession related, with less-pronounced modulation by eccentricity. Relatively higher sedimentation rates of the order of 3 to 5 cm kyr⁻¹ lead to pronounced precession cycle recordings, whereas slower sedimentation rates tend to amplify the modulation of precession cycles and thus eccentricity. Compared to sites with a lower sedimentation rate of 1 to 2 cm kyr⁻¹, the modulation of eccentricity is less pronounced in the XRF data of Site 1258. Both high-resolution bulk and benthic isotope records from Site 1258 revealed that multiple Eocene hyperthermal events exist, presumably forced by eccentricity-paced threshold passing (Sexton et al., 2011; Kirtland Turner et al., 2014). Recently, the eccentricity-driven variations in benthic $\delta^{13}\text{C}$ of sites 1262 and 1263, combined with the bulk $\delta^{13}\text{C}$ data from Site 1258, were used to construct an astrochronology by tuning the 405 and 100 kyr eccentricity components to the Laskar 2010d orbital solution (Laskar et al., 2011a; Lauretano et al., 2015, 2016). Due to a major shift in $\delta^{13}\text{C}$ at ~ 260 rmcd of Site 1263, two tuning options identifying two or three 405 kyr cycles in this interval were proposed. Integration of Site 1258 bulk isotope data and the best fit to the La2010d solution arguably led to a preferred age model with two 405 kyr cycles in the abovementioned interval at Site 1263 (Lauretano et al., 2016).

Based on the previous effort we construct a new consistent age model now spanning the entire Ypresian for sites from Leg 208 and Site 1258 by combining a wealth of available information with new high-resolution data, also making extensive use of the previously untouched spliced core images. Published benthic and bulk stable isotope data were combined for Leg 208 and Site 1258 (Figs. 2 and

S7), plotted at the Site 1263 revised meters composite depth (rmcd) and detrended for long-term trends (Fig. S9). Data have been linearly interpolated at 2 cm spacing and were then smoothed applying the Igor Pro smooth operation using binomial (Gaussian) smoothing and 30 001 points in the smoothing window. Benthic data from Site 1263 located in the disturbed drilling interval were removed from the combined record. The methods for time series analysis are those of Westerhold et al. (2015). Presence of the short and long eccentricity cycle in the isotope data is well documented (Zachos et al., 2010; Littler et al., 2014; Lauretano et al., 2015, 2016). Strong eccentricity-related cyclicity is clearly present in the evolutive wavelet power spectra of both isotope and XRF Fe data (Fig. S10) for the entire Ypresian, applying the magnetostratigraphic interpretation and using either the CK95 (Cande and Kent, 1995) or the GPTS2012 (Vandenberghe et al., 2012) ages for reversals.

We extracted the 405 and 100 kyr component of the data as detected in the evolutive spectra and plotted the filter over the data to investigate where the signal originates. The first-order tuning was done identifying the 405 kyr cycle in all data consistently. The advantage of having the high-resolution XRF Fe data is the ability to detect distinct modulation of the amplitude in the 100 kyr period related to the 405 and 2.4 myr eccentricity cycle modulations. In Fig. S11 in the Supplement, for example, around 242 and 265 rmcd of Site 1263 the XRF Fe data from Site 1263 and 1267 show very low amplitude variations separated by four 405 kyr cycles. If the amplitude modulation (AM) in the data is mainly driven by eccentricity (Zachos et al., 2010; Littler et al., 2014), then these intervals represent the 2.4 myr eccentricity nodes with minor amplitude variations on the 100 kyr level. Identification of the 2.4 myr minima is very important because they function as major tie points for orbital tuning and test for consistency with astronomical solutions (Westerhold et al., 2012, 2015; Zeeden et al., 2013). The starting point for the stable 405 kyr cyclostratigraphy is eccentricity cycle 119 at 48.0 Ma, which also represents a 2.4 myr eccentricity minimum (Fig. S11). Records presented here reconnect to astrochronologies that cover the Eocene cyclostratigraphic gap (Westerhold et al., 2015) from Site 1263.

4.1 How many 405 kyr cycles represent Chron C23?

Two enigmas had to be solved before a final stable 405 kyr cyclostratigraphy was set up. First, the question of whether two or three 405 kyr cycles are present at Site 1263 in the critical interval from 254 to 265 rmcd of Site 1263 needed to be answered. Second, the question of which orbital solution is appropriate for more detailed orbital tuning on the short eccentricity level needed to be answered. The first issue is complicated by two $\delta^{13}\text{C}$ shifts from 257 to 260 and at ~ 262 rmcd of Site 1263 (Fig. 2). Between the shifts, benthic $\delta^{13}\text{C}$ data at Site 1263 do not show a clear eccentricity-related cyclicity. Instead they reveal higher-frequency cycles

consistent with higher-frequency cycles in XRF Fe data of Sites 1258 and 1263. One argument to favor the model of two 405 kyr cycles is that it is consistent with uniform sedimentation rate above and below and across this interval (Lauretano et al., 2016). While a reasonable assumption in the absence of other constraints, sedimentation rates at Leg 208 sites do change across the ETM-2 event (see Littler et al., 2014, Fig. 7 therein), decreasing from 1.5 to 0.7 cm kyr⁻¹ at Site 1262. Using only the new magnetostratigraphy and the CK95 GPTS ages, sedimentation rates at Site 1263 drop from 2.6 to 2.0 cm kyr⁻¹ across the C24n.3n–C24r reversal close to ETM-2 and from 1.75 to 1.1 cm kyr⁻¹ across the C23n.2n–C23r reversal in the interval between the two $\delta^{13}\text{C}$ shifts (Fig. S13). A decrease in sedimentation rates at both Site 1258 and Site 1263 is very likely to be located in C23r. Time series analysis provides some evidence that cycle thicknesses change between 265 and 254 rmcd of Site 1263 (Fig. S11). Particularly relevant to this question is whether previous astrochronologies for Sites 1258 and 1263 (Westerhold and Röhl, 2009; Laurentano et al., 2016) were based on an underestimate of the duration of Chron C23n from Site 1258, which led to the 50 Ma discrepancy, thus making it difficult to determine the correct number of 405 kyr cycles in C23. Comparison to GPTS2012 (Vandenbergh et al., 2012) is compromised by the errors in the radioisotopic calibration points used in C22n (48.96 Ma \pm 0.33) and C24n (52.93 Ma \pm 0.23). Our new magnetostratigraphy is now more consistent with the width ratio observed in seafloor anomaly profiles from different ocean basins (Cande and Kent, 1995) that was used in GPTS2012.

The eccentricity-modulated precession cycles at ODP Site 1258 can help to test the effects of different numbers of 405 kyr cycles in the interval from 68 to 95 rmcd spanning Chron C23 (Fig. S14). The thicknesses of high-frequency cycles at Site 1258 change from \sim 45 to \sim 32 cm cycle⁻¹ across 77–82 rmcd at Site 1258 (Westerhold and Röhl, 2009, see Fig. S4 therein). Assuming that the average duration of the cycles is 21 kyr, the compression in cycle length translates into a decrease in sedimentation rate from 2.1 to 1.5 cm kyr⁻¹, as also seen in the new magnetostratigraphy. Applying the Laurentano et al. (2016) age models shows durations for the cycles of 16–24 kyr (two 405 kyr cycles) and 23–35 kyr (three 405 kyr cycles) in this interval (Fig. S14). Due to the constant sedimentation rates in their model between 68 and 92 rmcd, the shift in cycle thickness is transformed into overall shortening of the precession cycles, which is rather unrealistic. The option with two 405 kyr cycles seems to best fit the overall duration of 21 kyr for precession cycles and was thus chosen as the preferred age model by Laurentano et al. (2016) for Sites 1258 and 1263. We have developed a new 405 kyr age model (Table S46) based on the time series analysis of multiple high-resolution records that add new tie points in the interval from 68 to 92 rmcd previously not covered. Our new model proposed three rather than two 405 kyr cycles in this interval. As seen in the Site 1258 XRF Fe data,

the cycle thickness of the precession-related cycles for the entire interval now show a duration of 21–23 kyr (Fig. S14b). This basic age model also reveals that the long cycles of \sim 50 cm length around 80 rmcd represent 41 kyr obliquity cycles.

The principal terms for the precession of the Earth are given by the combination of the fundamental secular frequencies (g, s) of the solar system and the precession frequency p (Laskar, 1993; Hinnov, 2000; Laskar et al., 2004). Two periods dominate the precession: the \sim 23 kyr period related to Jupiter ($p + g5$) and Venus ($p + g2$) and the \sim 19 kyr period related to Mars ($p + g4$) and Earth ($p + g3$) (Laskar, 1999). Due to the evolution of the precession frequency p , the periods have been estimated to be \sim 22.5 and \sim 18.6 kyr (Laskar et al., 2004) 50 million years ago. In the 2.4 myr eccentricity minima, caused by the resonance between Earth and Mars ($g3 - g4$), the \sim 19 kyr period is weak or absent and only the \sim 23 kyr period is present (Laskar et al., 2004). It has been suggested that the amplitude modulation of the XRF Fe signal of Site 1258 spanning 80–87 rmcd could represent a 2.4 myr eccentricity minimum (Westerhold and Röhl, 2009). If this is correct, the dominating period for precession cycles recorded in the XRF Fe data for the interval from 68 to 92 rmcd should be towards 23 kyr rather than 19 kyr, rejecting the two 405 kyr cycle model (Lauretano et al., 2016). Because of a change in phasing between XRF Fe and bulk stable isotope data at 87 rmcd, cyclostratigraphy gets more complex at Site 1258. In contrast, Leg 208 sites show a consistent phase relation with decreased $\delta^{13}\text{C}$ values corresponding to higher Fe intensities. Both XRF Fe from Leg 208 and the combined carbon stable isotopes show three 405 kyr cycles in the 265–254 rmcd of Site 1263 interval (S11), and thus we propose a three-cycle model (Fig. S12).

4.2 Which orbital solution applies best for astronomical tuning?

4.2.1 Visual evaluation and determination

After a stable 405 kyr cyclostratigraphic framework is established, the orbital age model can be refined by tuning the carbon isotope data to an orbital solution on the short eccentricity level. Uncertainties in the ephemeris used to construct the orbital solutions currently limits their accuracy to roughly 48–50 Ma (Laskar et al., 2011a, b; Westerhold et al., 2015). Going beyond 50 Ma, the modulation pattern of short eccentricity recorded in the geological data can help to find the correct orbital solution (Laskar et al., 2004, 2011a, b). In particular, knowledge of the exact positions of the very long eccentricity minima, primary anchor for accurate orbital tuning (Shackleton and Crowhurst, 1997; Westerhold et al., 2007; Zeeden et al., 2014), can help to constrain astronomical solutions (Laskar et al., 2004; Westerhold et al., 2012). Beyond an age of 50 Ma, the position of the very long eccentricity nodes in available orbital solutions (La2004 – Laskar et al.,

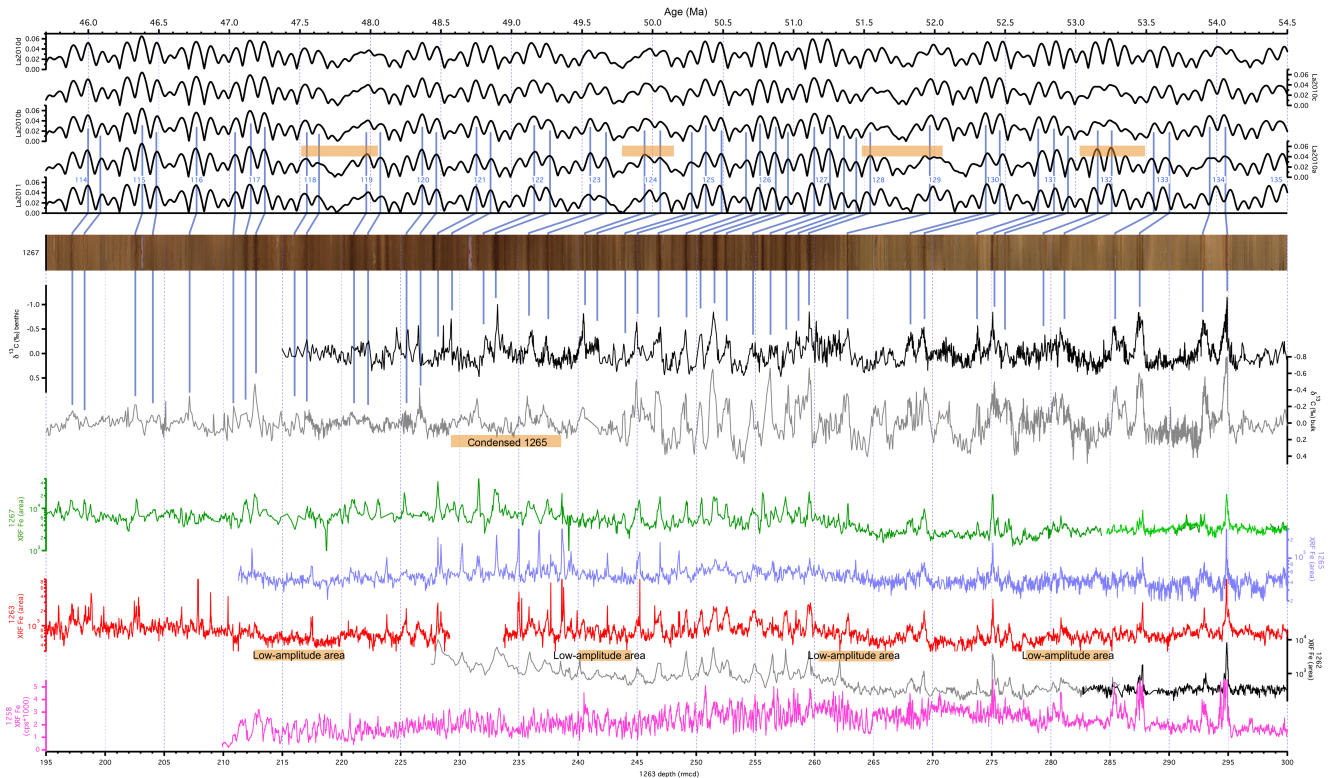


Figure 3. Correlation ties for astronomical calibration from 54.5 to 46 Ma. The top panel shows the numerical solutions La2011 (Laskar et al., 2011b) and La2010a-d (Laskar et al., 2011a), including the 405 kyr cycle number counted backwards from today (Wade and Pälike, 2004). Below the Site 1267 core image, which shows the best expressed dark layer pattern of all sites, we plotted the detrended benthic (black) and bulk (grey) combined $\delta^{13}\text{C}$ data as well as XRF Fe intensity records on revised meters composite depth from Site 1263. Areas of low-amplitude modulation in the data are marked in the XRF data for comparison. The blue lines show the tie points for the astronomically tuned age model correlating La2010b eccentricity maxima to lighter (more negative) $\delta^{13}\text{C}$ values. Note the condensed interval in the bulk Site 1265 $\delta^{13}\text{C}$ data, which cannot be tuned to the orbital solutions. For detailed discussion, see the text.

2004; La2010 – Laskar et al., 2011a; La2011 – Laskar et al., 2011b) are much more uncertain. Tuning to the La2010 or La2011 solution on the 100 kyr level is possible on the basis of a 405 kyr cyclostratigraphy but should be evaluated with great care. Before a bulk stable isotope record for Site 1258 and a benthic stable isotope record for Site 1263 were generated, only Site 1258 XRF Fe data provided a record sufficient for attempting astronomical tuning of the early Eocene (Westerhold and Röhl, 2009). Here we use a larger and more diverse data set from multiple sites on a stable 405 kyr cyclostratigraphy to test which orbital solution is the most appropriate one for detailed tuning.

The compiled records from the Ypresian, in particular the XRF core scanning Fe intensity data, show prominent minima in eccentricity-related modulation of the data in the intervals 212–220, 240–245, 260–267, 277–285, and 297–307 rmcd of Site 1263 (Figs. 3, S7). Starting with the first common node in the 2.4 myr cycle of La2010 and La2011 at 405 kyr cycles 118 to 119 (47.5–48.0 Ma), we go back in time and compare the position of the very long eccentricity cycle minima to the data amplitude minima of the Ypresian

records. Correlating the modulation minima at 212–220 rmcd of Site 1263 to the node at 405 kyr cycles 118 to 119 anchors the records to the astronomically tuned middle to late Eocene timescale (Westerhold et al., 2015). With the application of the stable 405 kyr framework introduced above, the preceding data modulation minima at 240–245, 260–267, 277–285, and 297–307 rmcd of Site 1263 require very long eccentricity minima at 405 kyr cycles 123–124, 128–129, 132, and 135–136 (Fig. 3). The first three minima can be observed in the different orbital solutions, suggesting that basically all the solutions back to 52 Ma could be used as target curves for tuning. Beyond 52 Ma, only the La2010b and La2010c solutions show a minimum at 405 kyr cycle 132. Going further back in time to 56–57 Ma, the minimum before ETM-2 (Lourens et al., 2005; Westerhold et al., 2007; Meyers, 2015) and the minimum before the PETM (Westerhold et al., 2007; Zachos et al., 2010; Littler et al., 2014) in the data even match very long eccentricity minima in La2010b and La2010c in 405 kyr cycles 135–136 (54.3–54.9 Ma) and 140 (56.4 Ma).

4.2.2 Statistical evaluation and determination

Extraction of the AM using statistical methods like those implemented in the astrochron package (Meyers, 2014) or the ENVELOPE (Schulz et al., 1999) routine are important for independently testing the visual recognition of cycle patterns (Hinnov, 2013; Hilgen et al., 2014). AM analysis on XRF core data using the ENVELOPE routine was applied at ODP Sites 1262 (52–60 Ma) and 1258 (47–54 Ma) records in Westerhold et al. (2012) in order to search for the very long eccentricity minima. Meyers (2015) used a^* values from ODP Site 1262 between PETM and ETM-2 to test the existing astrochronologies (Lourens et al., 2005; Westerhold et al., 2007). Both methods (Astrochron, ENVELOPE) thus provide sound statistical testing of chronologies at ODP Site 1258 and for Leg 208 sites.

Following the approach of Zeeden et al. (2015), we extracted the short eccentricity cycle (100 kyr) and applied a broad bandpass filter (0.004 to 0.016 cycles kyr⁻¹; 250–62.5 kyr cycle⁻¹; Tukey window) and subsequently made a Hilbert transform to extract the AM using the Astrochron software package (Meyers, 2015) for data from Sites 1258 and 1263. We applied the 405 kyr age model as a basic age model (Table S46). The resulting 405 kyr AMs of the XRF Fe intensity data are plotted against the La2004, La2010, and La2011 orbital solutions (Fig. 4). The AMs of the orbital solutions were extracted as described in Westerhold et al. (2012). For Site 1262 we plotted the 405 kyr AM of XRF Fe intensity data using the Option2 age model of Westerhold et al. (2012), which is almost identical to the updated 405 kyr age model presented here for the 53 to 58 Ma interval. We followed this procedure to demonstrate that similar results can be obtained with different approaches (Astrochron vs. ENVELOPE).

The position of the very long eccentricity minima in the AM of XRF Fe intensity data in the interval from 46 to 59 Ma (blue bars in Fig. 4) best fits with minima in the La2010b and La2010c orbital solutions. In contrast, the minima do not match minima in the La2004 solution, suggesting that this solution is not appropriate for testing geological data in this period of time. The La2010a–d and La2011 solutions fit geological data back to 50 Ma. Beyond 50 Ma these solutions diverge (as discussed in Westerhold et al., 2012). Only the La2010b and La2010c solutions exhibit very long eccentricity minima at ~53.3 and ~54.5 Ma. The minimum at ~54.5 Ma is a very prominent feature in the data of the Leg 208 sites that has been intensively discussed (Lourens et al., 2005; Westerhold et al., 2007; Meyers, 2015). The minimum at ~53.3 Ma is also detectable using the statistical methods but can be seen much better by visual inspection of the data (Figs. 2b, 3, and 5).

Quantitative evidence supporting the correct eccentricity node identification can also be derived from the emergence of obliquity cycles in the data at the nodes. Obliquity is not present in the Paleocene and early Eocene (Littler et al.,

2014; Zeebe et al., 2017) of the investigated records. However, Site 1258 Fe intensity data show some obliquity-related cycles at around 80 rmcd (also see Fig. S14) and from 55 to 60 rmcd corresponding to the end of the very long eccentricity node at 52 Ma and the beginning of the very long eccentricity node at 50 Ma. At another potential node (48 Ma), Site 1258 Fe data do not clearly exhibit obliquity cycles but rather low-amplitude modulations of precession-related cyclicity (Westerhold and Röhl, 2009, see Fig. 9 therein). Considering all these observations, they provide some independent evidence for the existence of eccentricity nodes at 50 and 52 Ma. The nodes at ~53.3 and ~54.5 Ma show no prominent obliquity cycles in the Fe records, as already discussed above and in Littler et al. (2014).

Based on our observations and the statistical analysis (Hinnov, 2013; Hilgen et al., 2014), we decided to fine-tune the records to the La2010b orbital solution (Fig. 5). It has to be noted here that there is hardly any difference between the La2010b and La2010c solution in the Ypresian from 46–56 Ma. Therefore, it does not matter if La2010b or La2010c is chosen as a target curve. Consequences of the match between orbital solutions and the geological data as well as the implications of the new age model for magnetostratigraphy need to be discussed. The tie points for the tuned short eccentricity age model are given in Table S47.

4.2.3 Potential distortion by nonlinear response of the climate system

Nonlinear response of climate is critical in the Ypresian. Multiple carbon cycle perturbations are documented as negative carbon isotope excursions (CIEs) and the dissolution of carbonates at the seafloor, both pointing to massive releases of ¹³C-depleted carbon to the ocean–atmosphere system. Associated warming has led to the term hyperthermals for these events. If the ¹³C-depleted carbon caused the warming and all events were triggered by carbon from the same reservoir, the magnitude of the CIEs scaled to the amount of carbon injected (Pagani et al., 2006). Because the residence time of carbon is of the order of 100 kyr (Broecker and Peng, 1982), the events will influence the amplitude of the bulk and benthic stable carbon isotope data and thus any AM analysis of early Eocene records. Additionally, the added carbon also leads to dissolution of carbonates at the seafloor, increasing the relative amount of non-carbonate material in the sediment (as detected by higher XRF Fe values). This will influence the statistical and visual recognition of cyclicity as discussed above. Modeling suggests that hyperthermals, except for the Paleocene–Eocene Thermal Maximum (PETM), could be paced by eccentricity forcing of the carbon cycle with the amplitudes of the events being partly driven by the eccentricity amplitude itself (Kirtland Turner et al., 2014). Variations in carbon isotope data do correspond to AM in the short-eccentricity cycle. In particular, the very long eccentricity minima are expressed as an interval of very low AM

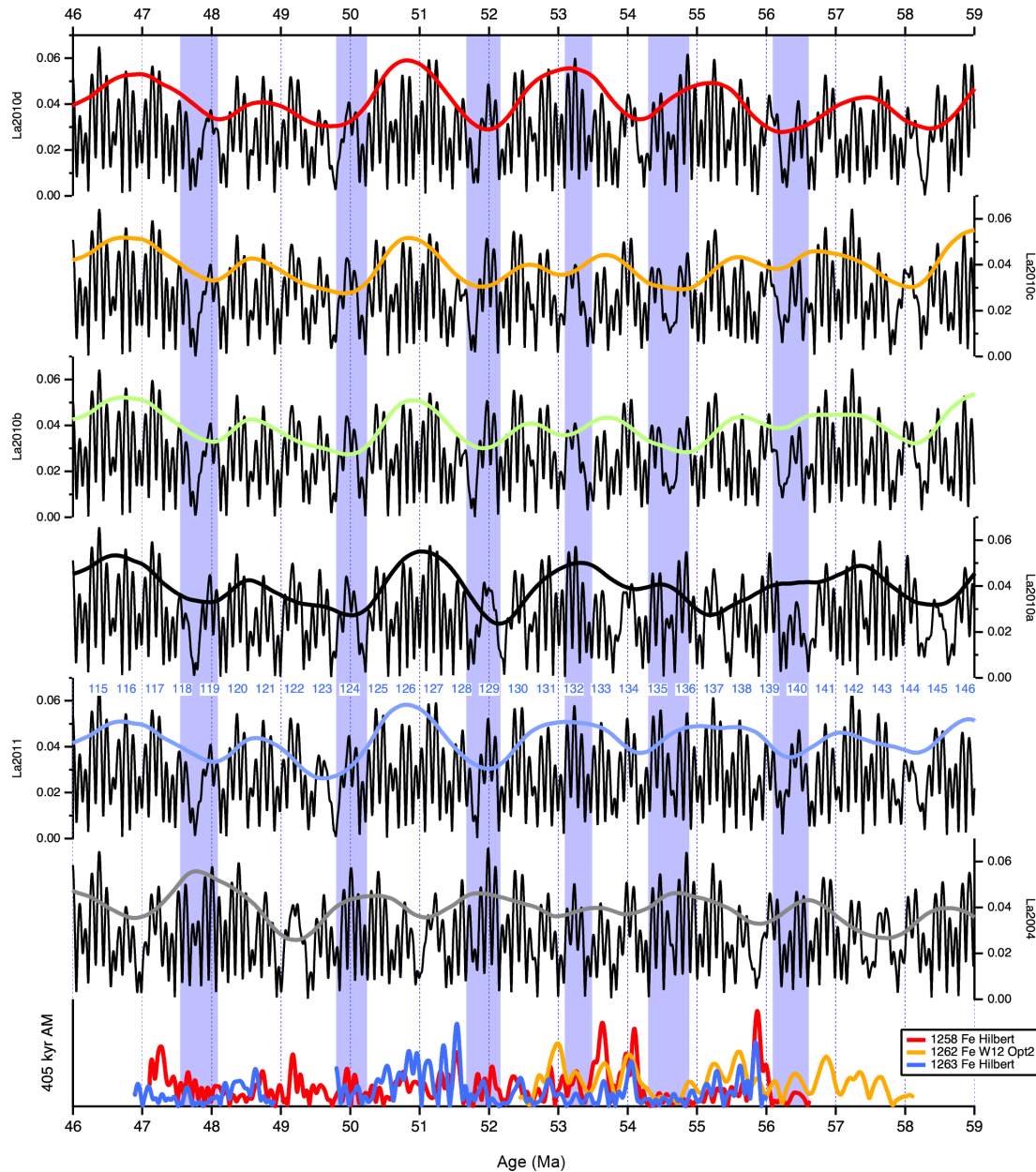


Figure 4. Comparison of the amplitude modulation (AM) of the short eccentricity cycle between the La2004, La2010, and La2011 orbital solutions and Fe intensity data from ODP Sites 1258 (red), 1262 (orange), and 1263 (blue). For the orbital solutions we also plotted the 405 kyr AM. The short eccentricity AM of Sites 1258, 1262, and 1263 Fe intensity data are plotted on the 405 kyr scale model. The very long eccentricity minima are highlighted by light blue bars in the orbital solutions and the Fe intensity data. Statistical and visual recognition of the cycle pattern suggests that the La2010b and La2010c solutions are most consistent with the geological data.

in the benthic carbon isotope data. Almost all hyperthermal events occur outside the very long eccentricity minima. Only very minor excursions at C21r5, C22r5, and C23n.2nH1 coincide with these nodes, but with comparatively reduced CIE than hyperthermals, suggesting that these might not be hyperthermals in the end. Hyperthermal layers are very well documented in the XRF data by prominent peaks due to dissolution of carbonate and since larger CIEs are characterized

by higher XRF Fe peaks. This tends to exaggerate the AM in the statistical analysis (Fig. 4). Hyperthermal events could be interpreted as amplifiers of the eccentricity amplitude, with a bias toward higher amplitudes. Because the focus in identifying the best fit astronomical solution lies on the very long eccentricity minima, this ensures that the distortion by hyperthermal events does not significantly alter the results of our study.

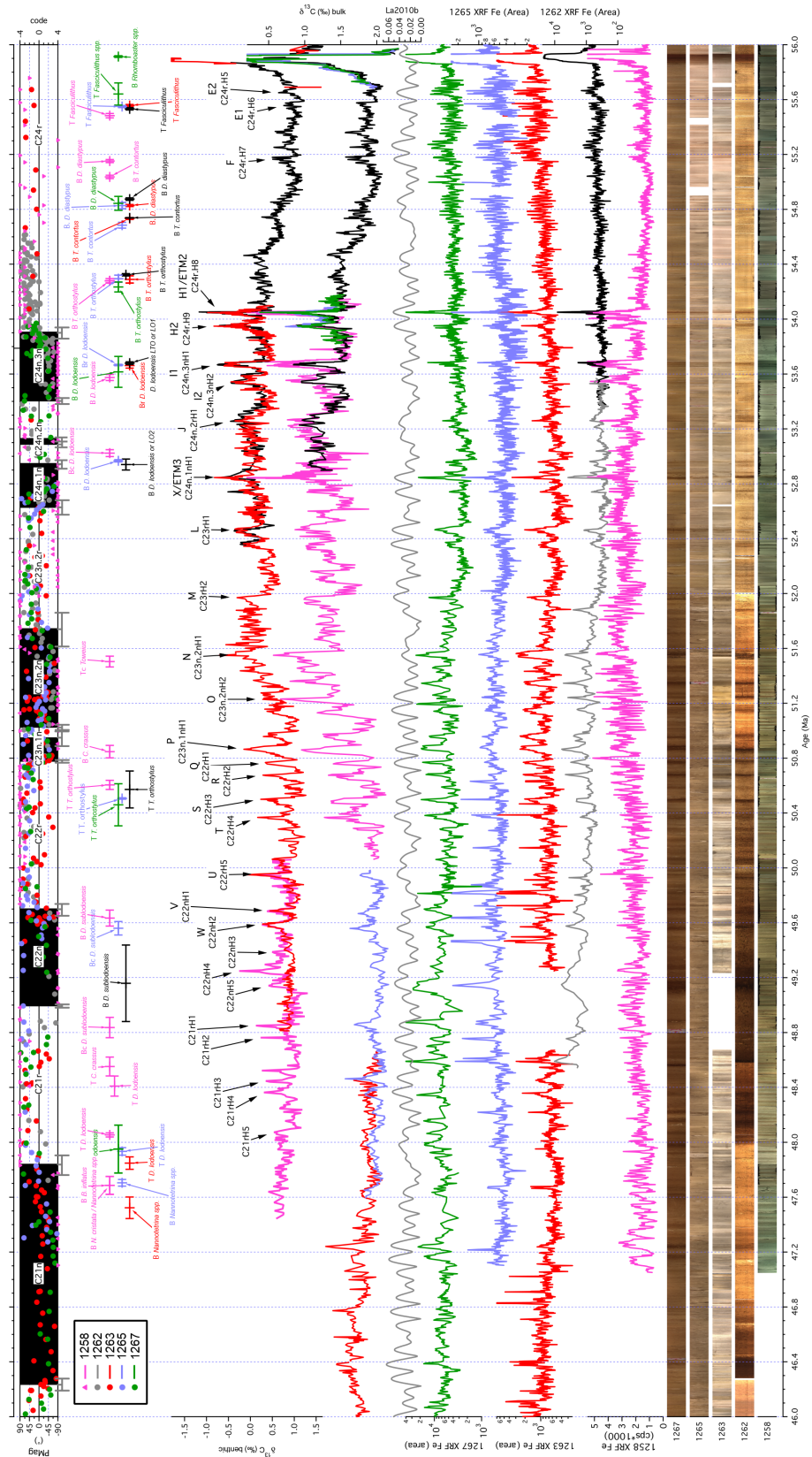


Figure 5. Bio- and magnetostratigraphic data, benthic and bulk stable isotopes, XRF core scanning Fe intensities, and core images for the new astrochronology for the Ypresian stage. Inclination data from ChRM analysis of each site from Leg 208 are shown. For Site 1258 the results of Suganuma and Ogg (2006) are given as a code between -4 and 4, with negative values indicating reversed polarity. The final synthesized magnetostratigraphy (Table S45) includes the error for the reversals (bars). Also shown is the La2010b (Laskar et al., 2011a) orbital solution by the central grey line. Note the extremely good match between orbital solution and amplitude modulation in the various XRF and stable isotope records. Hyperthermal events have been labeled according to Table 3.

5 Discussion

We have established the first complete astrochronology for the entire Ypresian stage (YATS), compiling, integrating, and synthesizing geochemical and bio- and magnetostratigraphic records at unprecedented precision. The result is a complex stratigraphy that can function as a reference, allowing synchronization of paleoclimate records essential to understanding cause and consequences of events in the early Eocene.

5.1 Geological evidence for chaotic behavior of the solar system in the Ypresian?

Just recently the first geologic evidence confirming the chaotic behavior of the solar system through the identification of a chaotic resonance transition during the Coniacian (~ 85 – 87 Ma) was reported (Ma et al., 2017). Similarly, the new records presented here appear to provide additional observational confirmation of the past chaotic evolution of the solar system, but as recently as ~ 52 Ma. Long-term simulations of orbital motions to study the stability of the solar system propose a chaotic rather than quasiperiodic pattern of motion in the solar system (Laskar, 1989; Laskar et al., 2004, 2011a). A macroscopic feature displaying the chaotic diffusion of the planetary orbits identifiable in geological records is the transition from libration to circulation in the resonant argument related to $\Theta = (s_4 - s_3) - 2(g_4 - g_3)$, the combination of angles in the precession motion of the orbits of Earth and Mars (Laskar et al., 2004, 2011a; Pälike et al., 2004; see Westerhold et al., 2015, chap. 5.3 for detailed discussion). One needs to extract the AM of both obliquity and precession in a geological data set in order to detect the transition from libration to circulation (Laskar, 1999), which is almost impossible, requiring a record that is influenced (or driven) by both high-latitude and low-latitude processes (Laskar et al., 2011a). Obliquity AM could be extracted from benthic $\delta_{18}\text{O}$ records, for example, if deep-sea temperature variations are continuously affected by obliquity. This is not the case for the Paleocene and early Eocene (Littler et al., 2014; Zeebe et al., 2017). Thus, investigation of the AM of obliquity is difficult with the currently available records. Laskar et al. (2011a) recommended searching for a modulation of the $g_4 - g_3$ period, the ~ 2.4 myr eccentricity modulation. The transition from libration to circulation should be visible by a switch from a ~ 2.4 myr period to a ~ 1.2 myr period in the modulation of eccentricity and climatic precession (Laskar, 1999; Laskar et al., 2004, 2011a; Pälike et al., 2004; Ma et al., 2017). Importantly, this could be transient with a switch back to ~ 2.4 myr shortly after (Laskar, 1999).

With the new cyclostratigraphy based on the stable 405 kyr eccentricity cycle for the Ypresian, we can test if this switch is present in the observations. Previously, the identification of the very long eccentricity cycle in geological data for the Ypresian used the XRF Fe intensities from Sites 1258 and 1262 only (Westerhold et al., 2007; Westerhold and

Röhl, 2009). The multiple proxy data now provide a much clearer picture, as described in Sect. 4.2, and show a very good match between geological data and the La2010b and La2010c numerical orbital solutions (Laskar et al., 2011a) for the Ypresian. The important feature shared between data and models is the position of the very long eccentricity minima expressed as areas of low-amplitude modulation in the data itself (Figs. 3, 4, 5). In the La2010b and La2010c solutions, the transition from libration to circulation occurs at ~ 52 Ma (Laskar et al., 2011a). Comparing orbital solutions with geological data indicates that the transition from libration to circulation occurred between 52 and 55 Ma. The AM minima in the data from 47 to 52 Ma are spaced at 2 to 2.4 myr (Figs. 4, 5), from 52 to 55 Ma they are spaced at roughly 1.2 myr, and after 55 Ma the spacing is ~ 2.4 myr.

None of the available orbital solutions perfectly fit the geological data. However, it is important that we isolated the transition in the data, which is also present in the La2010b and La2010c solutions. The short eccentricity cycle pattern both in the solutions and the geological data will not match perfectly beyond 50 Ma when the uncertainty in the solutions increase (as discussed in Westerhold et al., 2012). Still, the geological data and La2010b/c solutions are very similar from 53.5 Ma to the PETM. In the interval from 51 to 52 Ma, the most difficult part to tune in the Ypresian, multiple hyperthermal events and the shift in carbon isotope data make a direct comparison much more difficult. It has to be noted here that the eccentricity solutions from La2010b/c might not be completely reliable in this interval. Despite the uncertainties, we provide a tuned age model to La2010b/c because the match in the 53.5 Ma to the PETM interval is good enough to do so. If in doubt, the provided 405 kyr age model can still be used.

The point in time when the transition occurs in the numerical solutions is sensitive to the initial conditions of the planetary ephemeris used for back calculation of the planetary motions. The initial conditions depend on the accuracy of the observational data used to make a least squares fit of the model to the data. The La2010b/c solutions used the INPOP08 Ephemeris (Fienga et al., 2009; Laskar et al., 2011a). In contrast, the La2010d solution used the INPOP06 Ephemeris (Fienga et al., 2008) and the La2011 solution the INPOP10a ephemeris (Fienga et al., 2011). The very similar long-term behavior of INPOP06 and INPOP10a lead to the conclusion that these ephemerides are more stable than INPOP08 (Laskar et al., 2011b; Westerhold et al., 2012). Although the INPOP10a ephemeris is considered to be more accurate than INPOP08 (Fienga et al., 2011), the geological data provide evidence that the latter is closer to reality. Identifying the transition from libration to circulation at ~ 52 Ma in sediment archives is of great importance, not only because it supports the theory on the chaotic nature of the solar system (Laskar, 1989), but it also provides a benchmark to set the conditions for the gravitational model of the solar system (Laskar et al., 2004). It has to be noted that from an

Table 1. Comparison of magnetostratigraphic boundary ages in millions of years.

Chron	Standard GPTS			Tuned* GPTS			This study [†]		
	CK95	GPTS 2004	GPTS 2012	PEAT Sites [#]	Contessa Hyw	ODP Site 1260		ODP Site 1258 opt.2	ODP Site 1263
C18n.2n (o)	40.130	39.464	40.145	40.076 ± 0.005	40.120				
C19n (y)	41.257	40.439	41.154	41.075 ± 0.007	41.250	41.061 ± 0.009		41.030 ± 0.013	
C19n (o)	41.521	40.671	41.390	41.306 ± 0.005	41.510	41.261 ± 0.004		41.180 ± 0.011	
C20n (y)	42.536	41.590	42.301	42.188 ± 0.015	42.540	42.152 ± 0.007		42.107 ± 0.013	42.196 ± 0.013
C20n (o)	43.789	42.774	43.432		43.790	43.449 ± 0.018		43.517 ± 0.011	43.507 ± 0.011
C21n (y)	46.264	45.346	45.724		46.310			46.151 ± 0.009	46.235 ± 0.044
C21n (o)	47.906	47.235	47.349				47.723 ± 0.118	47.575 ± 0.018	47.834 ± 0.072
C22n (y)	49.037	48.599	48.566				48.954 ± 0.016		48.994 ± 0.012
C22n (o)	49.714	49.427	49.344				49.593 ± 0.042		49.695 ± 0.043
C23n.1n (y)	50.778	50.730	50.628				51.051 ± 0.021		50.777 ± 0.010
C23n.1n (o)	50.946	50.932	50.835				51.273 ± 0.039		50.942 ± 0.054
C23n.2n (y)	51.047	51.057	50.961				51.344 ± 0.032		51.025 ± 0.019
C23n.2n (o)	51.743	51.901	51.833				51.721 ± 0.023		51.737 ± 0.123
C24n.1n (y)	52.364	52.648	52.620				52.525 ± 0.023		52.628 ± 0.053
C24n.1n (o)	52.663	53.004	53.074				52.915 ± 0.029		52.941 ± 0.031
C24n.2n (y)	52.757	53.116	53.199				53.037		53.087 ± 0.021
C24n.2n (o)	52.801	53.167	53.274				53.111		53.123 ± 0.015
C24n.3n (y)	52.903	53.286	53.416				53.249 ± 0.017		53.403 ± 0.022
C24n.3n (o)	53.347	53.808	53.983				53.806 ± 0.020		53.899 ± 0.041

* Tuned GPTS compilation from Westerhold et al., (2015).

[†] Tuned to the orbital solution La2010b (Laskar et al., 2011a).

[#] Combined ages based on Pacific Equatorial Age Transect Sites 1218, U1333, and U1334 (Westerhold et al., 2014).

astronomer's point of view the La2010b and La2010c solutions are considered less reliable because they used a less stable ephemeris (Fienga et al., 2011; Laskar et al., 2011a). The same is true for the Laskar et al. (2004) solution, but the Niobrara data suggest a better fit to geological data than the nominal models of La2010d and La2011 (Ma et al., 2017). Clearly, the La2004 solution is not consistent with the geological data from 46 to 58 Ma (Fig. 4), with implications for arguments of Ma et al. (2017). In fact, all eccentricity AMs from the different theoretical astronomical solutions show a different and to some extent unusual behavior between 52 and 54 Ma (Fig. S15). Our new findings should motivate efforts to further explore the differences in the ephemerides.

Our results support the application of La2010b or La2010c solutions for eccentricity to construct astronomical age models back to 60 Ma. Beyond 60 Ma an accurate solution for eccentricity is not possible at the moment (Laskar et al., 2011b), but the stable 405 kyr cycles will still provide a good target to establish astrochronologies (Laskar et al., 2004, 2011a, b). Here again, the geological data should be examined to find the very long eccentricity minima in very early Cenozoic and Mesozoic strata (Meyers, 2015) and provide a landmark for developing more precise orbital solutions.

5.2 Solving the 50 Ma discrepancy in seafloor spreading rates

Combining the new astrochronology with the revised magnetostratigraphy for the Ypresian allows us to consider the significance of the abrupt global increase in spreading rates in Chron C23n.2n, which is also known as the 50 Ma dis-

crepancy on the Paleogene timescale (Vandenberghe et al., 2012). The unusual peak in spreading rates in the South Atlantic (Fig. 6a) is independent of the age model used for the magnetostratigraphic interpretation of Site 1258 (Westerhold and Röhl, 2009; Westerhold et al., 2012; Lauretano et al., 2016). The new multi-site magnetostratigraphic data from Leg 208 sites reveal that Chron C23n is too short in the Site 1258 magnetostratigraphic interpretation (Suganuma and Ogg, 2006) and the likely reason for the computed peak in spreading rates. Application of the tuned age model to the integrated magnetostratigraphy results in a moderate but distinct jump in spreading rates at the C23n.2n–C23r reversal from 12 to 19 km myr⁻¹. The exact age of the increase cannot be precisely located using the reversal pattern only, but it probably occurred somewhere in C23r or C23n.2n between 51.0 and 52.5 Ma. Interestingly, this timing is synchronous with a major reorganization of the plate–mantle system (Whittaker et al., 2007), the subduction initiation of the Izu–Bonin–Mariana arc (Ishizuka et al., 2011), and the bend in the Hawaii–Emperor seamount chain (O'Connor et al., 2013). Changes in spreading rates in the interval from 51.0 to 52.5 Ma thus seem to be a global phenomenon pointing to a major common driving mechanism.

Astronomical calibration of the refined magnetostratigraphy in the marine records resulted in an improved GPTS for the Ypresian (Fig. 6b, Table 1). Duration of polarity zones are now consistent with the reversal thickness relationships observed in the South Atlantic and within error for the mean width of magnetic anomalies as published in Table 4 of Cande and Kent (1992) (Tables 2 and S46). The improved magnetostratigraphy shows a 376 kyr shorter du-

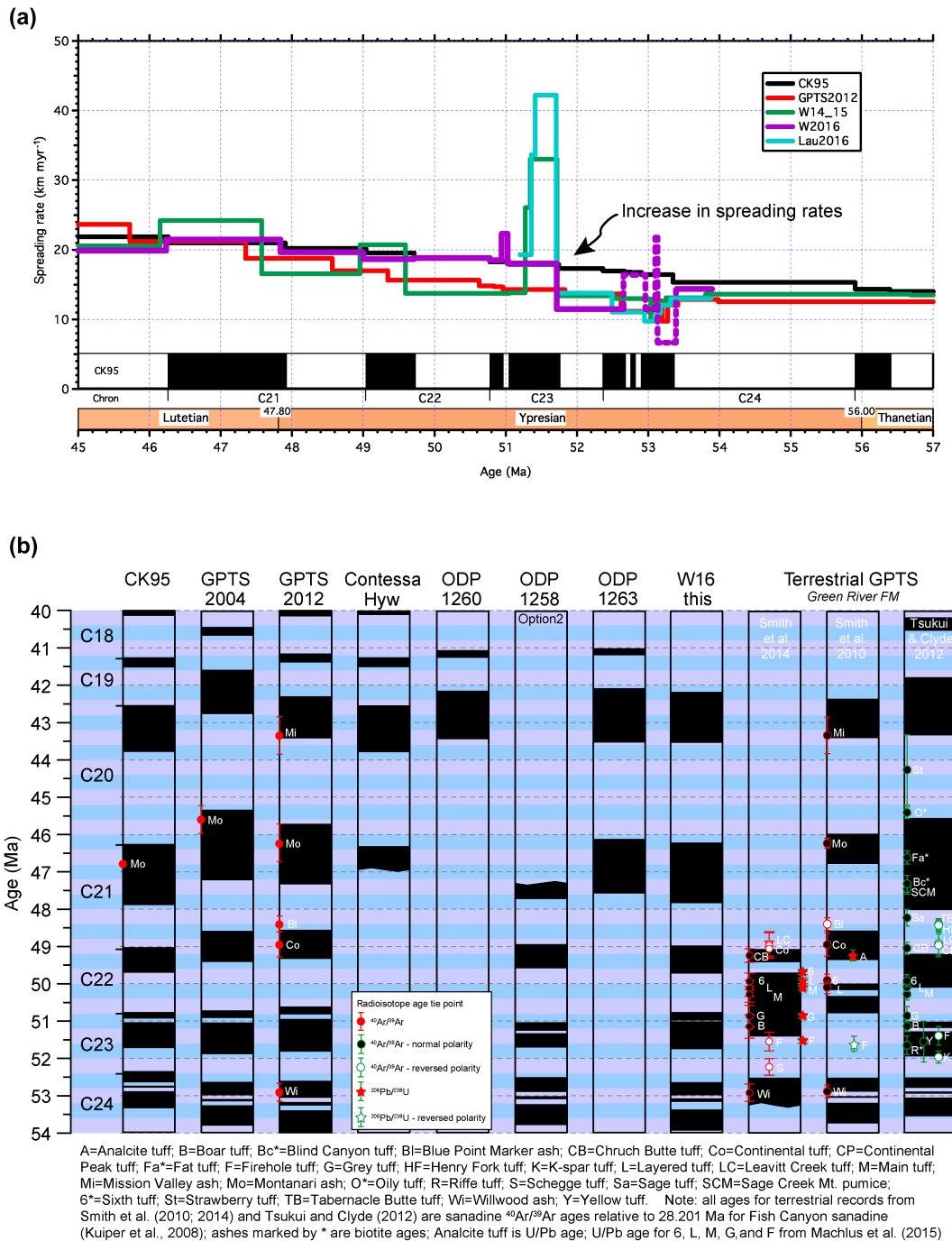


Figure 6. (a) Comparison of calculated spreading rates for the synthetic South Atlantic profile of Cande and Kent (1992) based on the age model of Cande and Kent (1995), GPTS2012 (Vandenberghé et al., 2012), Westerhold et al. (2015), Lauretano et al. (2016), and this study (W2016). The GPTS of CK95 is also plotted for reference. Note the increase in spreading rates in the South Atlantic at ~ 52 Ma based on the new GPTS presented in this study. (b) Geomagnetic polarity timescale of CK95 (Cande and Kent, 1995), GPTS2004 (Ogg and Smith, 2004), and GPTS2012 (Ogg, 2012; Vandenberghé et al., 2012) compared to astronomical calibrations of magnetochrons from Contessa Highway (Jovane et al., 2010), Site 1260 (Westerhold and Röhl, 2013), Site 1258 (Westerhold and Röhl, 2009; Westerhold et al., 2012), Site 1263 (Westerhold et al., 2015), and the new Ypresian GPTS (this study, W16) from 40 to 54 Ma. Terrestrial calibration of the GPTS from the Green River Formation (Smith et al., 2010, 2014; Tsukui and Clyde, 2012) is also shown. Small red dots with error bars mark the radiometric calibration points used for CK95, GPTS2004, GPTS2012, and Smith et al. (2010); green circles show calibration points for the terrestrial sections used by Tsukui and Clyde (2012). The overview demonstrates the now consistent Eocene GPTS from 30 to 54 Ma from ODP stratigraphic data and the discrepancy in the terrestrial GPTS. See text for discussion.

Table 2. Comparison of magnetostratigraphic boundary durations in millions of years.

Chron	Standard GPTS			Tuned* GPTS					This study [†] Leg208 Site1258
	CK95	GPTS 2004	GPTS 2012	PEAT Sites [#]	Contessa Hyw	ODP Site 1260	ODP Site 1258 opt.2	ODP Site 1263	
C18n.2r	1.127	0.975	1.009	0.999 ± 0.012					
C19n	0.264	0.232	0.236	0.231 ± 0.012		0.200 ± 0.007		0.150 ± 0.024	
C19r	1.015	0.919	0.911	0.882 ± 0.020		1.030		0.891 ± 0.006	
C20n	1.253	1.184	1.131			1.250		1.297 ± 0.013	
C20r	2.475	2.572	2.292			2.520		2.634 ± 0.020	
C21n	1.642	1.889	1.625					1.424 ± 0.027	
C21r	1.131	1.364	1.217					1.231 ± 0.134	
C22n	0.677	0.828	0.778					0.639 ± 0.058	
C22r	1.064	1.303	1.284					1.458 ± 0.063	
C23n.1n	0.168	0.202	0.207					0.222 ± 0.060	
C23n.1r	0.101	0.125	0.126					0.071 ± 0.071	
C23n.2n	0.696	0.844	0.872					0.377 ± 0.055	
C23n	0.965	1.205	1.200					0.670 ± 0.044	
C23r	0.621	0.747	0.787					0.804 ± 0.046	
C24n.1n	0.299	0.356	0.454					0.390 ± 0.052	
C24n.1r	0.094	0.112	0.125					0.122	
C24n.2n	0.044	0.051	0.075					0.074	
C24n.2r	0.102	0.119	0.142					0.138	
C24n.3n	0.444	0.522	0.567					0.557 ± 0.037	
C24n	0.983	1.363	1.360					1.554 ± 0.043	

* Tuned GPTS compilation from Westerhold et al., (2015).

[†] Tuned to the orbital solution La2010b (Laskar et al., 2011a).

[#] Combined ages based on Pacific Equatorial Age Transect Sites 1218, U1333, and U1334 (Westerhold et al., 2014).

ration for C22r, a 335 kyr longer duration for C23n.2n, and a 283 kyr shorter duration for C24n compared to previous marine records (Westerhold et al., 2015). The 50 Ma discrepancy in seafloor spreading rates is now eliminated. It was clearly the effect of the difficult and incomplete identification of Chron C23n at Site 1258. Moreover, the durations are consistent within error to the GPTS2012 (Vandenberghe et al., 2012) except for Chron C20r, which is difficult to access due to the relatively large error for the radioisotopic ages of the Mission Valley and the Montanari ash (Fig. 6) (for discussion see Westerhold et al., 2015). At this point more precise estimates of the mean width of magnetic anomalies and their error are required to be able to evaluate and improve the GPTS in the late Eocene. However, comparison to the GPTS models from terrestrial successions corroborates the finding (Westerhold et al., 2015) that the model of Tsukui and Clyde (2012) more closely resembles the marine GPTS than the model of Smith et al. (2010, 2014) (Fig. 6). Issues in the correlation of the Layered tuff, Sixth tuff, and Main tuff to local magnetostratigraphic records in the terrestrial records from the Green River Formation need to be resolved (see Tsukui and Clyde, 2012) to understand the current discrepancies in various terrestrial and marine GPTS models.

5.3 Defining the age of the top and bottom of the Ypresian stage

The GSSP of the Ypresian, which also marks the Paleocene–Eocene boundary, is defined at the basal inflection of the CIE of the PETM (Aubry et al., 2007), about two-thirds of the

way down in magnetostratigraphic C24r (Westerhold et al., 2007) at the base of Zone CNE1 where the top of the calcareous nannofossil *Fasciculithus richardii* group and the base of calcareous nannofossil excursion taxa (CNET) occur (Westerhold et al., 2007, 2015; Agnini et al., 2014). The age and position of the onset of the PETM is confirmed by our study and thus needs no further discussion.

The top of the Ypresian stage (or base of the Lutetian) is defined at the lowest occurrence of the calcareous nannofossil *Blackites inflatus* (CP12a–CP12b boundary; Okada and Bukry, 1980) in the Gorrondatxe sea cliff section in Basque Country, northern Spain (Molina et al., 2011). The lowest occurrence of *B. inflatus* is reported ~ 819 kyr (39 precession cycles) after the base of Chron C21r (Bernaola et al., 2006; Payros et al., 2009), or 60 % up in Chron C21r (C21r.6), leading to an age of ~ 47.8 Ma using GPTS2012 (Vandenberghe et al., 2012). The new YATS absolute ages for the C21n–C21r (47.834 Ma) and C21r–C22n (48.994 Ma) boundaries result in a 500 kyr older top of the Ypresian (= base of *B. inflatus*; C21r.6 will be 48.3 Ma). However, when data from the Gorrondatxe reference section are compared with those from Site 1258 and SE Newfoundland Ridge (Norris et al., 2014), a large discrepancy emerges in the position of the base of *B. inflatus*, which is located near the bottom of Chron C21n outside the GSSP section (Fig. 5) and has a tuned age of 47.686 ± 0.065 Ma. According to Agnini et al. (2014), *B. inflatus* is more abundant in shallow water or hemipelagic settings and is rarely observed in pelagic settings. Furthermore, and importantly, few data exist that provide a precise tie to magnetostratigraphy. Diachronicity for some nannofos-

Table 3. Overview of naming and age of hyperthermal events.

Hyperthermal naming scheme				Tuned age Ma La2010b	Comment
Sexton et al., (2011)	Kirtland-Turner et al., (2014)	Cramer et al., (2003) Lauretano et al., (2016)	This study		
C21rH6	C21rH5	-	C21rH5	48.075	
C21rH5	C21rH4	-	C21rH4	48.365	
C21rH4	C21rH3	-	C21rH3	48.450	
C21rH3	-	-	-	-	not crossing the 1 σ level of Kirtland-Turner et al., 2014
C21rH2	C21rH2	-	C21rH2	48.765	
C21rH1	C21rH1	-	C21rH1	48.850	
C22nH4	C22nH5	-	C22nH5	49.140	
C22nH3	C22nH4	-	C22nH4	49.250	
-	C22nH3	-	C22nH3	49.385	added by Kirtland-Turner et al., 2014
C22nH2	C22nH2	W	C22nH2	49.585	
C22nH1	C22nH1	V	C22nH1	49.685	
C22rH3	C22rH6	-	-	-	no clear d13C and d18O excursion
C22rH2	C22rH5	U	C22rH5	49.950	
C22rH1	-	-	-	-	not crossing the 1 σ level of Kirtland-Turner et al., 2014
-	C22rH4	T	C22rH4	50.370	
-	C22rH3	S	C22rH3	50.485	
-	C22rH2	R	C22rH2	50.670	
-	C22rH1	Q	C22rH1	50.760	
-	C23n.1nH1	P	C23n.1nH1	50.865	
-	C23n.2nH1	O	C23n.2nH2	51.230	name adjusted to revised magnetostratigraphy
-	C23rH3	N	C23n.2nH1	51.550	name adjusted to revised magnetostratigraphy
-	C23rH2	M	C23rH2	51.970	
-	C23rH1	L	C23rH1	52.460	
-	C24n.1nH1	K/ETM-3	C24n.1nH1	52.845	
-	-	J	C24n.2rH1	53.260	not crossing the 1 σ level of Kirtland-Turner et al., 2014
-	C24n.3nH2	I2	C24n.3nH2	53.545	
-	C24n.3nH1	I1	C24n.3nH1	53.665	
-	C24rH9	H2	C24rH9	53.950	
-	C24rH8	H1/ETM-2	C24rH8	54.050	
-	C24rH7	F	C24rH7	55.165	
-	C24rH6	E2	C24rH6	55.555	
-	C24rH5	E1	C24rH5	55.650	
-	PETM	PETM	PETM	55.930	

sil events is well documented in the high-resolution astronomically calibrated bio- and magnetostratigraphy for Leg 208 sites and Site 1258 presented here. However, an offset of more than 600 kyr (48.3–47.686 Ma) between the Integrated Ocean Drilling Program (IODP) Atlantic sites and the Gorrondatxe section, resulting from the inconsistent position in C21n rather than C21r.6, suggests that either the base of *B. inflatus* might not be the perfect global marker for the Ypresian–Lutetian boundary or the quality of magnetostratigraphic data in the Spanish section is not appropriate. Additional work on deep marine and hemipelagic records with good magnetostratigraphy is required to test the consistency of the position of this biohorizon and the relative ranking and spacing with the base of *Nannotetrina cristata* that is reported to occur very close to the base of *B. inflatus* and Chron C21n (Agnini et al., 2014; Norris et al., 2014).

5.4 Hyperthermal events

The Ypresian stage is of special interest to recent scientific work because it allows the study of climate dynamics and feedbacks in a warm world (Zachos et al., 2008). In particular, the occurrence of multiple transient global warming events (hyperthermals) could help tremendously to understand the response of the climate system to a massive release of carbon to the ocean–atmosphere system (Dickens, 2003; Zachos et al., 2008; Lunt et al., 2011; Kirtland Turner et al., 2014). A paired negative excursion in the carbon and oxygen isotope composition of bulk sediment and benthic foraminifera associated with a more clay-rich layer, indicating dissolution of carbonate, are the characteristics of the deep-marine hyperthermal events (Zachos et al., 2005; Lourens et al., 2005; Leon-Rodriguez and Dickens, 2010). The early Eocene hyperthermals are paced by Earth’s orbital eccentricity, except for the PETM (Zachos et al., 2010; Sexton et al., 2011; Littler et al., 2014; Lauretano et al., 2015, 2016; Laurin et al., 2016). After the discovery of a large number of early Eocene hyperthermal events, a magnetostratigraphic

based naming scheme was introduced (Sexton et al., 2011; Kirtland Turner et al., 2014). Because the scheme used the inaccurate magnetostratigraphy of Site 1258 for Chrons C21 to C24, we updated the naming scheme of Kirtland Turner et al. (2014). We also maintain the labeling system of Cramer et al. (2003), which was extended by Lauretano et al. (2016) for consistency. In Table 3 we provide an overview of the naming schemes, astronomical age, and position with respect to magnetostratigraphy.

The most important characteristic used to identify hyperthermal events is a paired excursion in $\delta^{13}\text{C}$ and $\delta^{18}\text{O}$ isotope data of preferably benthic foraminifera as in Sexton et al. (2011). However, some of the events defined in Sexton et al. (2011) fail this criterion but have still been considered as hyperthermals if the excursions are more negative than 1σ beyond the mean of its 2 Myr running average (Kirtland Turner et al., 2014). This kind of test was applied to both benthic and bulk stable isotope data. It can be questioned if the bulk stable isotope data alone are sufficient to identify a hyperthermal event as several other factors can influence bulk sediment composition (e.g., production and preservation). To definitely identify hyperthermal events in the entire Ypresian, paired high-resolution benthic stable $\delta^{13}\text{C}$ and $\delta^{18}\text{O}$ values are needed. These records are available from PETM to the C24n.2rH1 event (Littler et al., 2014) and from the C22rH5 to C21rH5 events (Sexton et al., 2011). Unfortunately, the published record from Site 1263 (Lauretano et al., 2016), which spans from ETM-2 (H1) to the C22nH2 (W) event, only provides benthic $\delta^{13}\text{C}$ data. To definitely identify hyperthermals in the Ypresian, more stable paired benthic isotope data are needed. Hence, we refrain from changing or even making up a new naming scheme. We synthesized the available schemes and adopted the naming in Table 3 here as well as provided their astronomical ages from the new astrochronology (Fig. 5).

6 Conclusions

A new complex cyclostratigraphy and refined bio-, chemo-, and magnetostratigraphy have been developed for key ODP records spanning the entire Ypresian stage from 56 to 47 Ma. Detailed correlation of ODP Sites 1258, 1262, 1263, 1265, and 1267 using the new CODD macros software tool revealed a 3–400 kyr condensed interval at the Leg 208 sites during Chron C22n. New characteristic remanent magnetization data from four Leg 208 sites show an overall consistent magnetostratigraphy refining the Ypresian geomagnetic polarity timescale. Multi-site ChRM data correlated on a centimeter to decimeter scale suggest that a magnetostratigraphic record from a single site might contain significant errors due to coring disturbance. Cyclic variations in synthesized XRF core scanning and stable isotope data as well as lithological changes apparent in core images have been successfully used to refine previous astrochronologies and con-

struct the first complete Ypresian astronomical timescale. In the absence of independent, high-precision time-control-like radioisotopic dates in Leg 207 and 208 sediments, our study clearly validates that it is crucial to combine multiple records from multiple regions to help safeguard against incompleteness that is otherwise difficult to assess both qualitatively and quantitatively. The YATS not only provides updated absolute ages for bio- and magnetostratigraphy but also a comprehensive list of the early Eocene hyperthermal events. The new astronomically calibrated Ypresian GPTS resolves the 50 Ma discrepancy, which was primarily caused by the imprecise magnetostratigraphy of Site 1258. Comparing the eccentricity-related cyclic pattern in XRF core scanning and stable carbon isotope data to numerical orbital solutions suggests that the transition from libration to circulation as predicted by the La2010b solution occurred ~ 52 Ma. This adds to the geological evidence for the chaotic nature in the evolution of the solar system.

Data availability. The data reported in this paper are tabulated in the Supplement and archived in the PANGAEA (<https://doi.pangaea.de/10.1594/PANGAEA.871246>; Westerhold et al., 2017) database.

The Supplement related to this article is available online at <https://doi.org/10.5194/cp-13-1129-2017-supplement>.

Competing interests. The authors declare that they have no conflict of interest.

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