

Hidden in Plain Sight: The Overlooked Influence of the Cs⁺ Substructure on Transformations in Cesium Lead Halide Nanocrystals

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Cesium lead halide perovskite nanocrystals, with formula CsPbX₃ (X = Cl, Br, I), have received enormous attention in the last five years because of their efficient, narrow, and tunable light emission.¹ They crystallize in structures characterized by a three-dimensional network of corner-sharing [PbX₆]⁴⁻ octahedra (either the perfectly cubic α , or the slightly distorted orthorhombic γ and tetragonal β structures). However, perovskites are only some of the many compositions and structures found in the Cs–Pb–X ternary system: together with the two binary compounds CsX and PbX₂, the best-known structures in the system are Cs₄PbX₆, the perovskites α, β, γ -CsPbX₃, the nonperovskite δ -CsPbX₃, and CsPb₂X₅. The corresponding nanocrystals are highly reactive systems that can readily undergo chemical and structural transformations in response to changes in their surroundings.^{1,2} Rationalizing the reactivity of these materials is crucial for two reasons. First, for their stabilization. Second, to create opportunities for processability and applications.

The reports describing compositional and structural interconversions are often focused on pairs of compounds, e.g., the CsPbX₃ \rightleftharpoons Cs₄PbX₆ or the CsPbX₃ \rightleftharpoons Cs₄Pb₂X₅. However, as more experimental data have become available over time, it is now possible to connect different reactions and construct a unifying picture of the reactivity of the Cs–Pb–X compounds. Such a view has been missing to date, mainly because these materials are often approached with a focus on how their structures are related to their properties and less to their reactivity. In fact, [PbX₆]⁴⁻ octahedra are usually considered as the main structural motif in lead halide perovskite and perovskite-related structures. Such octahedra are disconnected in Cs₄PbX₆, but share corners in the perovskite CsPbX₃ and faces in the δ -CsPbX₃ nonperovskite structure. This rationalization is useful when discussing electronic properties, as they mostly depend on the Pb–X–Pb connectivity.³ However, reasoning in terms of [PbX₆]⁴⁻ octahedra excludes many of the materials in the Cs–Pb–X system that can either generate or be generated from one of the three above through a chemical reaction and hides deeper structural connections across the whole series. For example, CsX nanocrystals can be transformed into CsPbX₃ nanocrystals,⁴ and CsPb₂X₅ can be obtained by thermal

decomposition of CsPbX₃ nanocrystals.⁵ Neither CsX nor CsPb₂X₅ actually contains [PbX₆]⁴⁻ octahedra: lead in CsPb₂X₅ features an unusual [PbX₈]⁶⁻ biaugmented triangular prism octa-coordination.⁶

In this Viewpoint, we present a different way of approaching this whole set of compounds. Instead of [PbX₆]⁴⁻ octahedra, we take the arrangement of Cs⁺ cations as the leading structural motif (Figure 1). From such a perspective, the

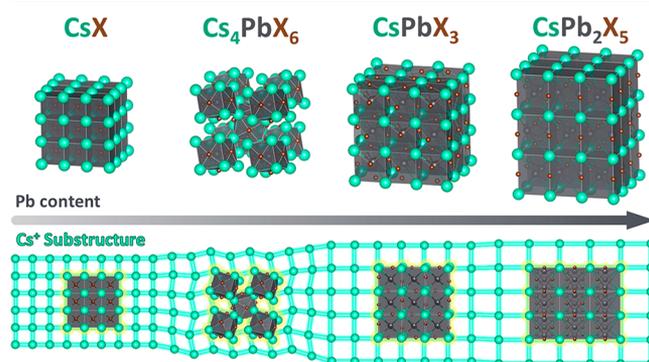


Figure 1. Crystal structures of CsX, Cs₄PbX₆, CsPbX₃, and CsPb₂X₅ compounds shown in order of increasing lead content (top panel) and interconnected by a common structural motif: a substructure of Cs⁺ cations that expands and distorts along with the series (bottom panel).

parent structure of the group is CsX, where the Cs⁺ ions form a substructure of cubic cages, each enclosing a single halide ion. Such a substructure is found in all the other compounds (except for δ -CsPbX₃), albeit expanded and sometimes distorted. Starting from the parent structure CsX, we can replace X⁻ ions with composite [Pb_{*n*}X_{2*n*+1}]⁻ polyanions, which

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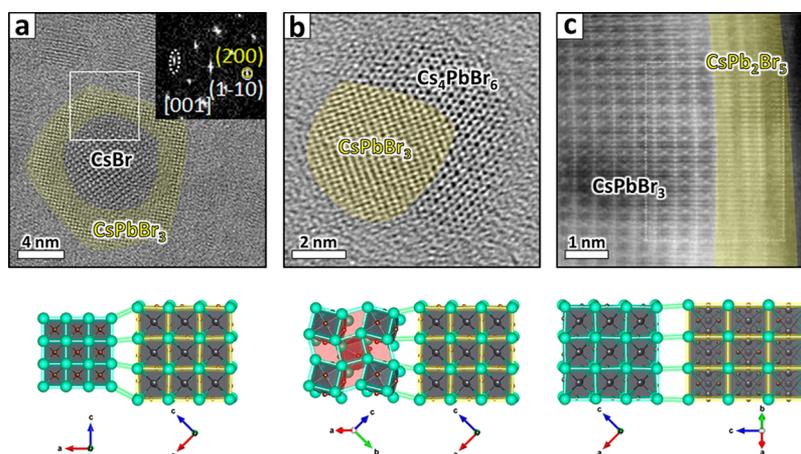


Figure 2. Epitaxial heterostructures between pairs of Cs–Pb–Br compounds. From left to right: high-resolution TEM images of CsBr–CsPbBr₃, Cs₄PbBr₆–CsPbBr₃ and CsPbBr₃–CsPb₂Br₅ heterostructures.^{4,7,8} The bottom sketches are simplified models of the epitaxial connections, highlighting the expansion and distortion of the Cs⁺ substructure when moving from lead-poor to lead-rich compounds. The lead-rich compounds are shaded in yellow, and the empty cages in the Cs₄PbBr₆ structure are shaded in red. Crystallographic axes are shown to clarify the relative orientation of the crystal structures. The HRTEM image in panel a is reproduced from ref 7; those in panels b and c are reproduced from refs 4 and 8 with permission from the Royal Society of Chemistry and Springer Nature, respectively.

is equivalent to adding n units of PbX₂ to each unit of CsX. Depending on the amount of PbX₂ added, the anionic network adapts its connectivity to ensure the charge balance, and we obtain the whole series of structures.

We begin with Cs₄PbX₆, where only one-fourth of the Cs⁺ cages incorporate lead in the form of disconnected [PbX₆]⁴⁻ ions, resulting in an average [Pb_{0.25}X_{1.5}]⁻ polyanion per cage. Note that the filled cages retain a cubic shape, while the empty ones distort to become rhombohedral (Figure 2, the empty cages in the Cs₄PbBr₆ structure are shaded in red). That distortion makes the Cs⁺ cations in the corner of one cube interact with X⁻ anions on the faces of the surrounding cubes and produces an overall change in symmetry from cubic to hexagonal. In CsPbX₃, all the cages become filled with [PbX₃]⁻. In each cage, the center is taken by Pb²⁺ ions, while X⁻ ions are on the faces, shared between neighboring cages. This is equivalent to the more familiar description of the structure through corner-sharing [PbX₆]⁴⁻ octahedra. Moving along the series, in CsPb₂X₅, each cage encloses a [Pb₂X₅]⁻ polyanion. Here, the two Pb²⁺ cations and four of the five X⁻ anions are inside the cage, while the remaining X⁻ anion is found at one edge and is therefore shared with the neighboring cages. This configuration leads to a connected network of [PbX₈]⁶⁻ biaugmented triangular prisms. PbBr₂, which has no Cs⁺ substructure, should not be considered as next in the list. However, it is interesting to notice that the Pb²⁺ ions in it share the same coordination found in CsPb₂X₅, although the connectivity changes to keep the structure charge-balanced. In this sense, PbX₂ can be imagined as an infinitely large [Pb_{*n*}X_{2*n*+1}]⁻ network (with $n \rightarrow \infty$) enclosed in a Cs⁺ cage where the Cs–Cs distance has expanded to infinity.

The existence of a common structural motif helps to explain the reactivity of cesium lead halide compounds: in essence, the preservation of the Cs⁺ substructure allows the transformation to take place without tearing the whole crystal structure apart. In our discussion, we will be focusing mainly on the Cs–Pb–Br system, because the many reports on these materials provide enough experimental data to test our mechanistic hypothesis. Especially, we will refer to nanocrystals since they are ideal systems to study transformations because

of their fast reaction kinetics and the possibility of direct observations with transmission electron microscopy (TEM). However, the structural relationships between compounds exist regardless of the crystal size, thus it applies to bulk materials as well.

The simplest structure in the series is CsBr. Shamsi et al. achieved the transformation of CsBr nanocrystals to γ -CsPbBr₃ perovskite nanocrystals by reaction with Pb(oleate)₂.⁴ The transformation preserved the nanocrystal size distribution, pointing against a complete dissolution–recrystallization mechanism (i.e., in which all CsBr is dissolved and CsPbBr₃ is nucleated from free ions in solution). Instead, the reaction proceeded by the progressive insertion of PbBr₂ into the CsBr structure, with Pb²⁺ coming from Pb(oleate)₂ and Br⁻ from a partial etching of the nanocrystal surface. The reported observation of intermediate CsBr@CsPbBr₃ core–shell epitaxial heterostructures supports our proposed mechanism (Figure 2a). The average size of the heterostructures became smaller as the transformation proceeded, accounting for the partial surface etching and the release of excess Cs⁺ in the environment. At the same time, the epitaxial connection ensured that the Cs⁺ substructure was shared between the two domains. As expected, the large difference of crystal structure constants (CsBr = 4.3 Å; CsPbBr₃ = 5.8 Å) produced strain at the interface, which was measured by high-resolution TEM.

The next and much more extensively studied set of transformations concerns the conversion of CsPbBr₃ to Cs₄PbBr₆ and back. This equilibrium was reported to be reversible and easily triggerable by a variety of stimuli: addition and subtraction of PbBr₂ or CsBr and changes in the concentration of ligands in the environment. One interesting example was reported by Li et al., who cycled the transformation back and forth while retaining the nanocrystal size distribution.⁹ Specifically, the CsPbBr₃ → Cs₄PbBr₆ reaction was triggered by adding a mixture of ligands (oleic acid and oleylamine), which increased the PbBr₂ solubility in the environment, while the inverse Cs₄PbBr₆ → CsPbBr₃ reaction was driven by the direct addition of PbBr₂. Remarkably, they observed that the nanocrystal size at each cycle of the transformation was compatible with a constant number of Cs⁺

atoms. All these observations are consistent with a substructure-preserving reaction mechanism: PbBr_2 is readily incorporated in (or released by) each Cs^+ cage, as the process requires only a mild distortion of the Cs^+ substructure. Baranov et al. provided a strong proof for this mechanism: while studying the $\text{Cs}_4\text{PbBr}_6 \rightarrow \text{CsPbBr}_3$ transformation, they isolated intermediate $\text{CsPbBr}_3\text{-Cs}_4\text{PbBr}_6$ epitaxial heterostructures where the Cs^+ substructure was continuous across the whole nanocrystals, with little distortion at the interfacial regions (Figure 2b).⁷ Interestingly, the transformation proceeded from one single nucleation point on the nanocrystal surface, and from there it extended to the remaining regions. The absence of multiple nucleation points of the CsPbBr_3 phase on a single nanocrystal can be explained based on the proposed mechanism: once the transformation is started, the strain induced by the epitaxial connection locally straightens the distorted Cs^+ substructure of Cs_4PbBr_6 to match that of CsPbBr_3 . This local expansion facilitates the further diffusion of PbBr_2 at the domain interface, resulting in a sort of self-catalyzed reaction.

Another interesting example is the $\text{Cs}_4\text{PbBr}_6 \rightarrow \text{CsPbBr}_3$ transformation driven by CsBr extraction. This was reported upon the reaction of nanocrystals with solid Prussian Blue ($\text{Fe}^{\text{III}}_4[\text{Fe}^{\text{II}}(\text{CN})_6]_3$),¹⁰ which is known to selectively intercalate Cs^+ , or after their exposure to a hexane–water interface.¹¹ In both cases, no intermediate heterostructures were isolated, although the size distribution of the nanocrystals was preserved, evidence that again played against a dissolution–recrystallization mechanism. Also, the nanocrystals contracted in size after the reaction, as expected from a partial etching. This is a key point to rationalize the observations based on our proposed transformation mechanism. Cs_4PbBr_6 nanocrystals indeed should undergo etching when CsBr is extracted from their surface. However, for each four CsBr units etched away, one unit of PbBr_2 is also freed. No PbBr_2 was detected at the end of the process by XRD, and that can be explained by PbBr_2 re-entering the remaining portion of the Cs_4PbBr_6 nanocrystals, driving the transformation to CsPbBr_3 nanocrystals. The extraction of Cs^+ ions stopped when what remained of the etched nanocrystals had been entirely converted to CsPbBr_3 : any further extraction of CsBr would have been accompanied by the extraction of PbBr_2 , which, as already said, was not detected in the final product. Furthermore, Wu et al. compared the volume of particles before and after the water-driven transformation, finding it compatible with a constant number of Pb atoms (Cs_4PbBr_6 diameter = 17.8 nm; CsPbBr_3 expected/measured = 10.8/12.2 nm).¹¹ We repeated the analysis for the Prussian Blue driven transformation, finding an excellent match (Cs_4PbBr_6 diameter = 9.8 nm; CsPbBr_3 expected/measured = 5.9/6.3 nm). This further supports our mechanistic hypothesis. It is worth mentioning that a dissolution–recrystallization mechanism in these types of reactions is still possible, given the right conditions. For example, Palazon et al. and Udayabhaskararao et al. both triggered the $\text{CsPbX}_3 \rightarrow \text{Cs}_4\text{PbX}_6$ transformation by adding amines, which can sequester lead by forming a complex.^{12,13} However, ammonium ions are known to compete with Cs^+ ions in the lead halide perovskite structure, thus causing the Cs^+ substructure degradation and the consequent recrystallization.^{14,15}

To conclude the series, fewer reports exist on the $\text{CsPbBr}_3 \rightarrow \text{CsPb}_2\text{Br}_5$ transformation, mainly because of the less

appealing optical and electronic properties of CsPb_2Br_5 . In general, this transformation is slower than the previous ones or requires harsher conditions such as continuous illumination or heating.^{2,16} For example, Huang et al. reacted CsPbBr_3 nanosheets with an excess of PbBr_2 at 60 °C for 3 days, obtaining epitaxial heterostructures.¹⁷ Because there is little interest in pure CsPb_2Br_5 , the transformation was halted at that stage. The preservation of the Cs^+ substructure helps to rationalize these observations as well. The $\text{CsPbBr}_3 \rightarrow \text{CsPb}_2\text{Br}_5$ transformation stretches the Cs^+ cages in one direction, lowering their symmetry from cubic to tetragonal. As a consequence, the crystal structure parameter in the elongated direction becomes too large (9.7 Å) to match with that of CsPbBr_3 , leaving only the (001) plane as a possible interface between CsPbBr_3 and CsPb_2Br_5 . The high-quality HRTEM images captured by Zheng et al. on a cosynthesized heterostructure help to visualize this connection (Figure 2c).⁸ The main consequence of this structural constraint is that the transformation cannot proceed cell-by-cell as for the previous examples: it must take place layer by layer, largely increasing the required activation energy and slowing down the reaction.

Overall, the preservation of the Cs^+ substructure weaves a thread connecting the various reactions in the Cs-Pb-X system. Its main merit is to outline a common reaction mechanism, based on the intercalation/extraction of PbBr_2 from lead-poor/rich compounds, which proceeds via epitaxial heterostructures and preserves the nanocrystal size and size distribution. In this sense, trying to predict or model the response of the Cs^+ substructure to transformations can offer new insights into those reactions that we expect to occur but have not been observed to date. For example, the $\text{CsBr} \rightarrow \text{CsPbBr}_3$ transformation seems to bypass the anticipated Cs_4PbBr_6 intermediate. The reason is likely found in the self-catalyzed $\text{Cs}_4\text{PbBr}_6 \rightarrow \text{CsPbBr}_3$ transformation discussed above: Cs_4PbBr_6 might form in the early stages of the reaction, but the incoming PbBr_2 will likely react more favorably with it, forming CsPbBr_3 , than with CsBr . Yet, small domains of Cs_4PbBr_6 , although not observed to date, may exist at the CsBr-CsPbBr_3 interface, acting as a strain-releasing transient layer.

Moving to other halides, an interesting observation is that both CsI and Cs_4PbI_6 nanocrystals react with $\text{Pb}(\text{oleate})_2$ and PbI_2 , respectively, forming the metastable but desirable black CsPbI_3 perovskite phase instead of the thermodynamically favored wide-gap $\delta\text{-CsPbI}_3$.^{4,18} The preservation of the Cs^+ substructure helps to rationalize that reaction route. The arrangement of Cs^+ ions in $\delta\text{-CsPbI}_3$ is completely different from that found in CsI and Cs_4PbI_6 ; thus, the transformation from CsI or Cs_4PbI_6 to $\delta\text{-CsPbI}_3$ would require a major structural rearrangement of the Cs^+ substructure and hence high activation energy. Instead, both CsI and Cs_4PbI_6 transform to the more structurally compatible CsPbI_3 perovskite structure, both transformations requiring much lower activation energy. In nanocrystals, the perovskite $\text{CsPbI}_3 \rightarrow \delta\text{-CsPbI}_3$ transformation slowly happens soon after, likely by a recrystallization mechanism, posing a challenge to the applications of this promising material, for example, in solar cells.^{19,20}

Another example of a process preserving the Cs^+ substructure is the $\gamma\text{-CsPbBr}_3 \rightarrow \gamma\text{-CsPbI}_3$ halide-exchange reaction.^{21,22} The reverse transformation in bulk compounds has been recently reported for the nonperovskite phases, $\delta\text{-CsPbI}_3 \rightarrow \delta\text{-CsPbBr}_3$.²³ These two reactions involve

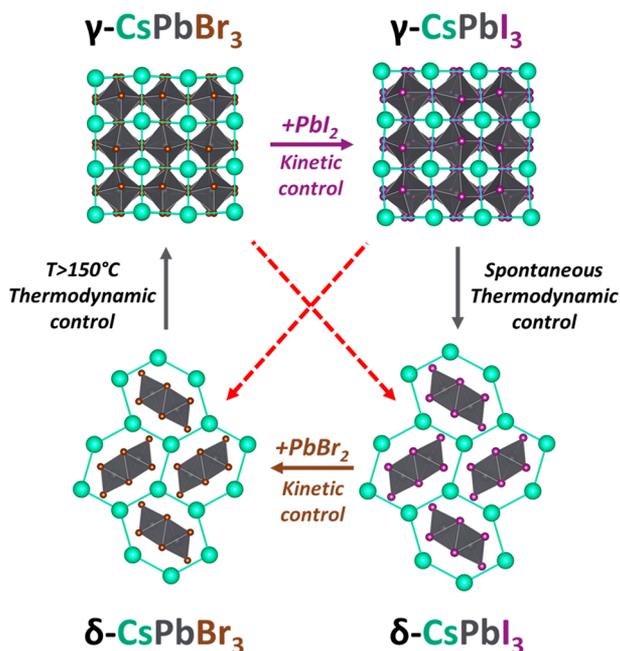


Figure 3. Summary of the transformations between the γ -perovskite and δ -nonperovskite polymorphs of CsPbBr_3 and CsPbI_3 . Ion-exchange reactions (indicated by purple and brown arrows) that preserve the Cs^+ substructure are under kinetic control and produce metastable structures. Instead, the thermodynamically favored structures are generated by reactions (indicated by red dashed arrows) that do not preserve the Cs^+ substructure. The metastable structures convert to the stable ones either spontaneously in nanocrystals (CsPbI_3) or upon heating in bulk powders (CsPbBr_3).

completely different structures yet share similarities (Figure 3). Both transform a stable structure (γ - CsPbBr_3 or δ - CsPbI_3) into a metastable one (γ - CsPbI_3 or δ - CsPbBr_3 , respectively) under kinetic control, and both preserve the Cs^+ substructure. The Cs^+ substructure acts as a backbone in both γ and δ structures (the arrangement of $[\text{PbX}_6]^{4-}$ octahedra is preserved as well) during anion exchange. The further transformation to the stable γ - CsPbBr_3 phase could be achieved only by raising the temperature above 150 °C, indicating the high activation energy required to disrupt and rearrange the Cs^+ substructure.

The last case we discuss is that of Cs_2PbX_4 . This stoichiometry is intermediate between Cs_4PbX_6 and CsPbX_3 , corresponds to Cs^+ cages filled with $[\text{Pb}_{0.5}\text{X}_2]^-$ anions on average, and exists in the form of the Ruddlesden–Popper $\text{Cs}_2\text{PbCl}_2\text{I}_2$ structure (Figure 4, middle model).²⁴ Here, the difference in size between Cl^- and I^- helps to stabilize the structure. The $\text{Cs}_2\text{PbCl}_2\text{I}_2$ structure features cubic Cs^+ cages filled with $[\text{PbCl}_2\text{I}_2]^{2-}$ alternated with distorted noncubic empty cages, for an average $[\text{Pb}_{0.5}\text{ClI}]^-$ unit per cage, and hence fits into the structural series discussed so far. The filled cages are shifted by half-cell to let Cs^+ in one layer interact with I^- in the neighboring one and provide a stabilizing interaction for the structure, similarly to Cs_4PbX_6 . Reactions involving Ruddlesden–Popper metal halide nanocrystals are rare; however, a recent work reported the transformation of $\text{Cs}_2\text{PbCl}_2\text{I}_2$ nanoplates into nanocrystals of γ - CsPbBr_3 and Cs_4PbCl_6 by reaction with PbBr_2 and MnCl_2 , respectively.²⁵ Although there is not sufficient data to exclude recrystallization, it is worth pointing out that the Cs^+ substructure in $\text{Cs}_2\text{PbCl}_2\text{I}_2$ can be derived from that of γ - CsPbX_3 by

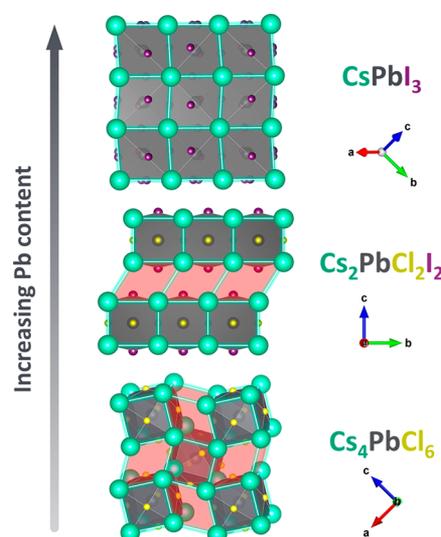


Figure 4. Ruddlesden–Popper metal halide $\text{Cs}_2\text{PbCl}_2\text{I}_2$ (middle) has an intermediate structure between that of lead-rich CsPbI_3 (top) and the lead-poor Cs_4PbCl_6 (bottom), into which it can be transformed. The three structures share similarities in the cubic Cs^+ substructures, albeit with different numbers of cubic cages filled by lead halide polyanions. Empty cages in the front layer are highlighted in red. Crystallographic axes are shown to clarify the relative orientation of the Cs^+ substructure and the unit cell.

alternating distortions of one layer of cages (the empty ones) while leaving the other layer almost intact (the filled ones).

The outcomes of the $\text{Cs}_2\text{PbCl}_2\text{I}_2$ reactions with PbBr_2 and MnCl_2 can be explained through the principle of Cs^+ substructure preservation. The addition of MnCl_2 would cause an $\text{I} \rightarrow \text{Cl}$ halide exchange to a hypothetical Cs_2PbCl_4 structure that is not stable and must decompose. In parallel, Mn^{2+} is not able to replace Pb^{2+} in CsPbCl_3 in large amounts,²⁶ leaving the material the only option of expelling PbI_2 and converting into stable Cs_4PbCl_6 . From the Cs^+ substructure point of view, this requires tilting of the cages, as half of them become empty. Instead, the addition of PbBr_2 proceeds with the replacement of both I^- and Cl^- by an intermediate-sized Br^- along with further incorporation of PbBr_2 , which removes the distortions of the Cs^+ substructure and delivers a stable γ - CsPbBr_3 . Such analysis yields a prediction that lead halide nanocrystals of Ruddlesden–Popper compounds could be obtained from Cs_4PbX_6 by PbX_2 insertion. That line of thinking opens a perspective onto Cs_4PbX_6 nanocrystals as precursors for multinary metal halide nanocrystals through the reactive intercalation of PbX_2 and other metal halide M_nX_m species.

To conclude, the discussed preservation of the Cs^+ substructure could be an example of the general “A” cation substructure preservation applicable to other halides. For example, compounds where Cs^+ is replaced by organic cations such as methylammonium or formamidinium often crystallize in analogous structures (e.g., α, β, γ - MAPbX_3 and MA_4PbX_6 ,^{27,28} where MA = methylammonium). Studies on their reactivity are highly focused on the substructure-destructive γ - $\text{MAPbI}_3 \rightarrow \delta$ - MAPbI_3 transformation that poses huge challenges to the development of perovskite-based photovoltaics.²⁹ However, we expect organic cation-based halides to show substructure-conservative transforma-

tions like those of cesium compounds. The “A” cation substructure preservation is also expected to hold in structures where the B cation is different from Pb^{2+} , such as Sn^{2+} or Ge^{2+} . Further on, it will be interesting to explore if a similar principle applies to more complex structures such as the double perovskites, where the pairs of Pb^{2+} cations of CsPbX_3 are replaced by B^+ and B^{3+} cations, opening up to investigation a vastly untapped array of possible transformations.

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Notes

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